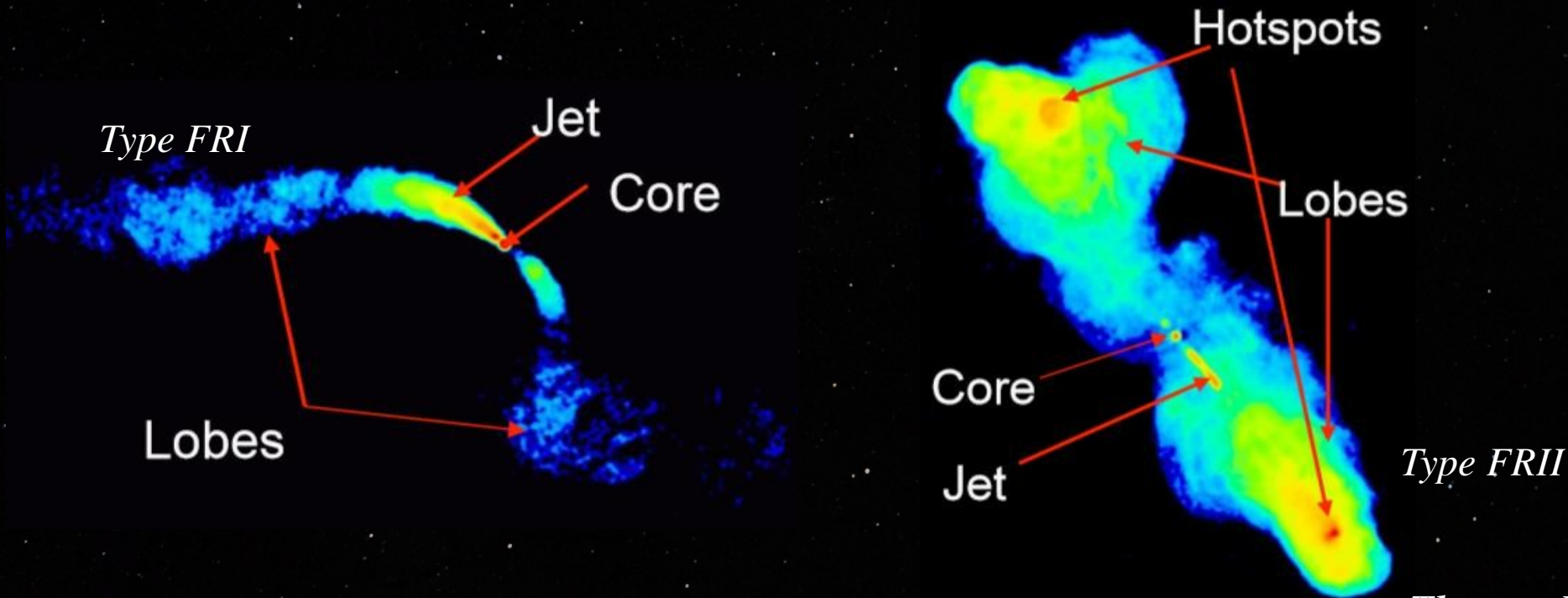


Jet asymptotic collimation : how does it depend on the source properties ?



Thomas Jannaud (IPAG, Grenoble)

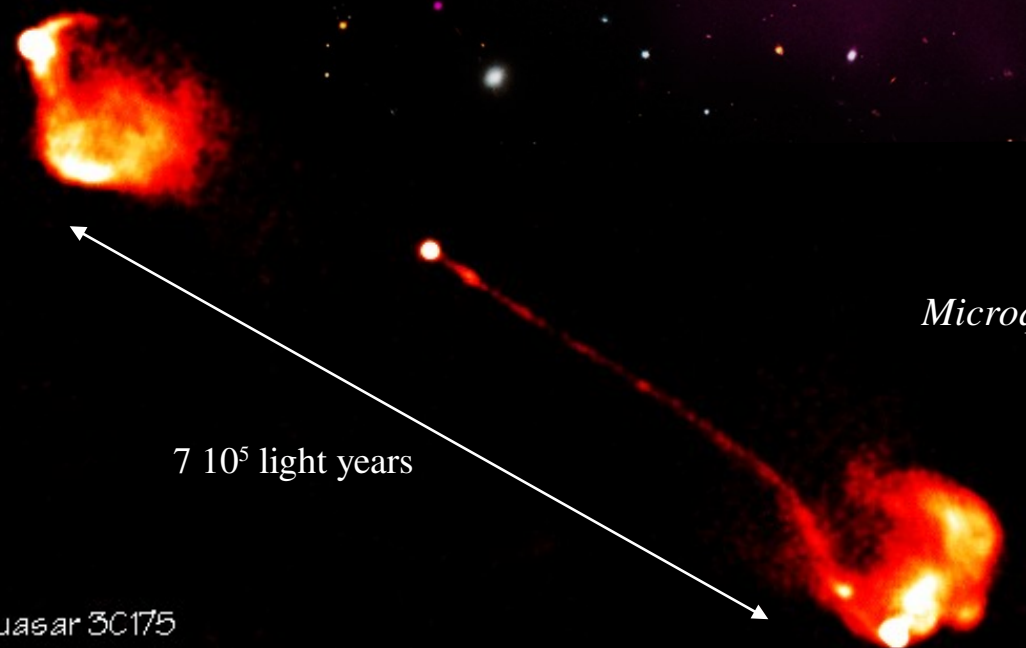
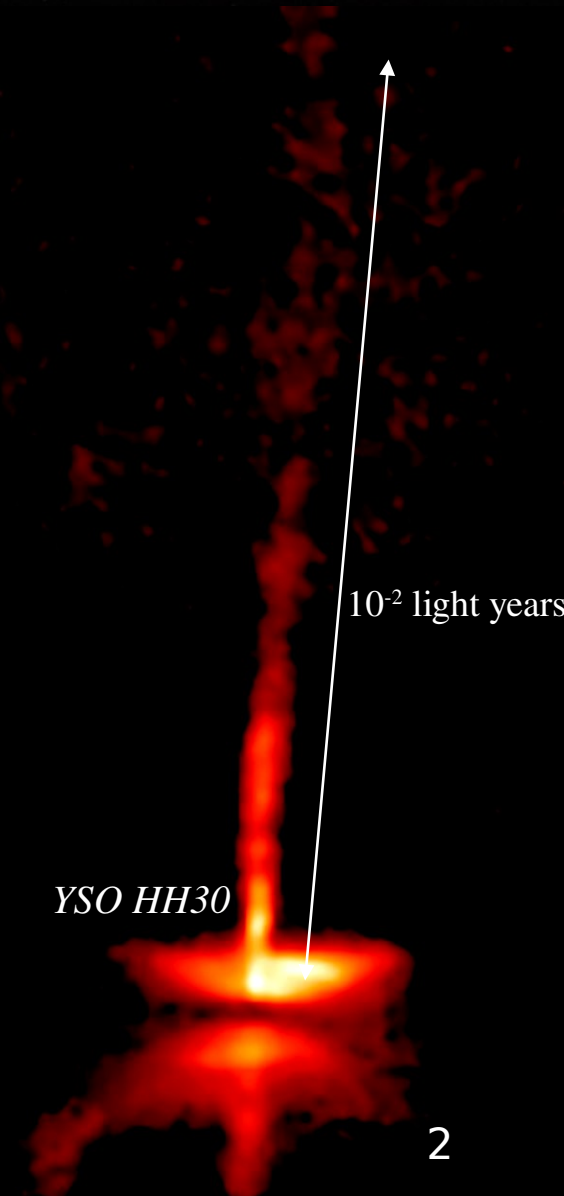
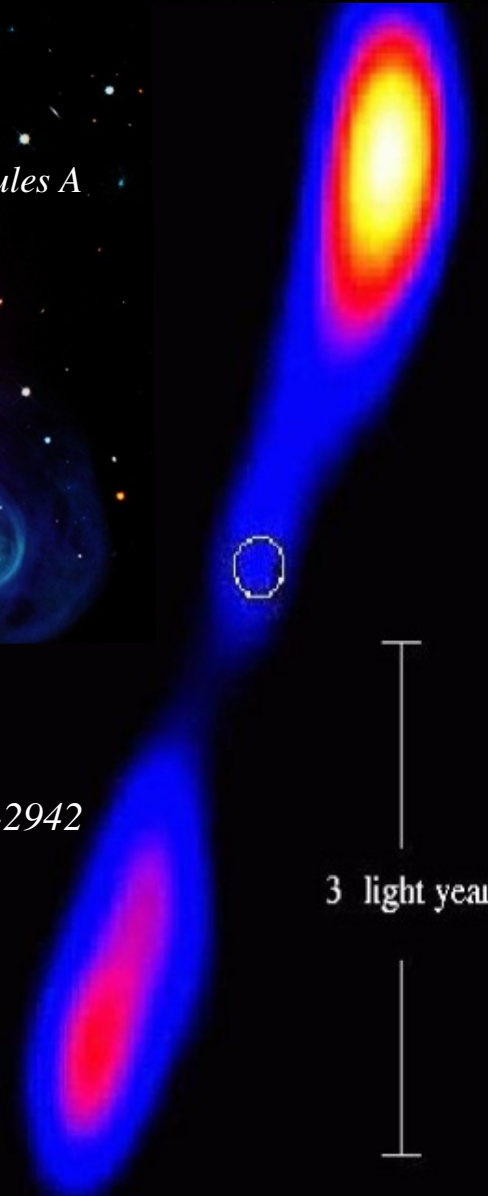
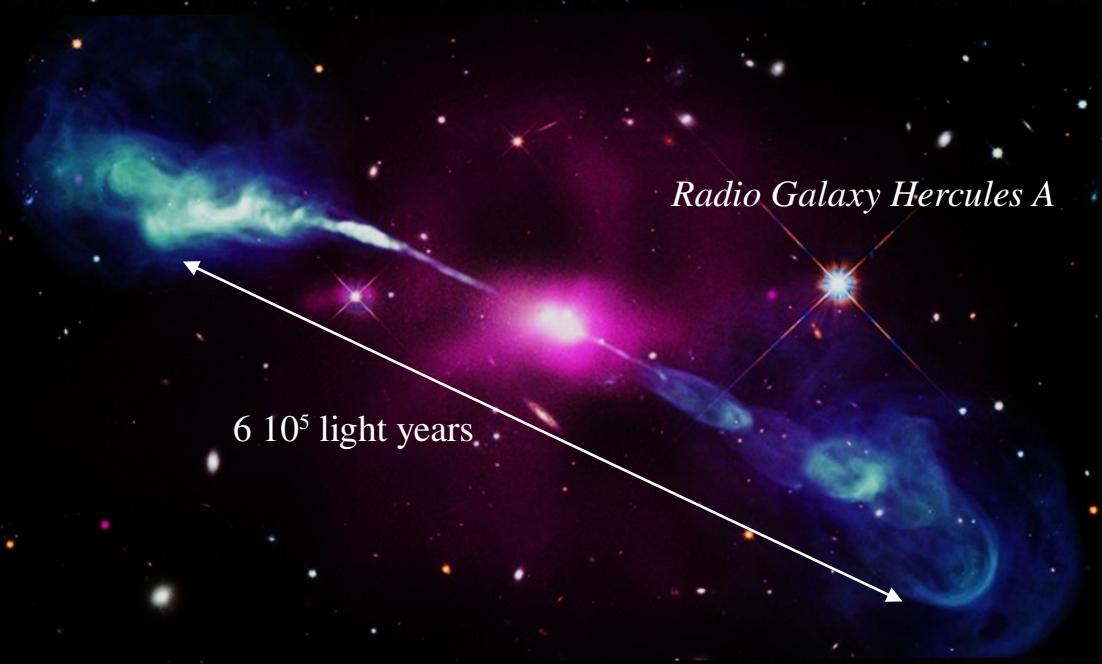
Jonathan Ferreira (IPAG, Grenoble)

Claudio Zanni (INAF, Turin)

March 31, 2022

Presentation for the 10th FER0 meeting

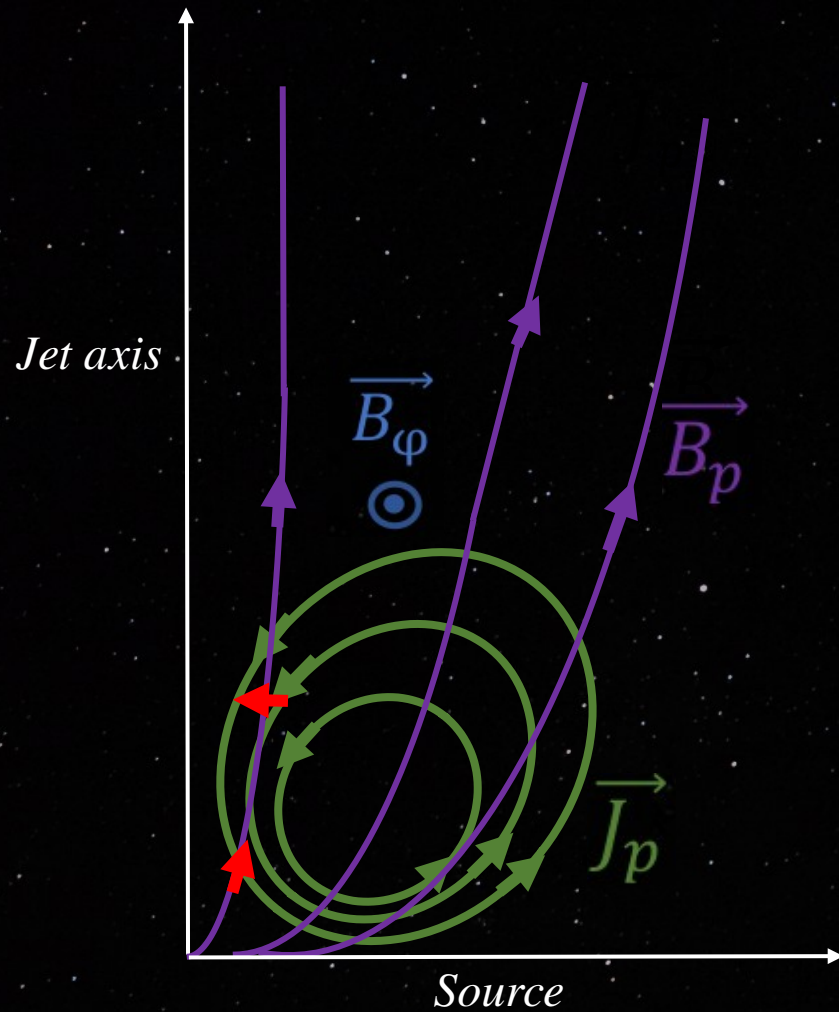
Astrophysical jets



Quasar 3C175
YLA 6cm image (c) NRAO 1996

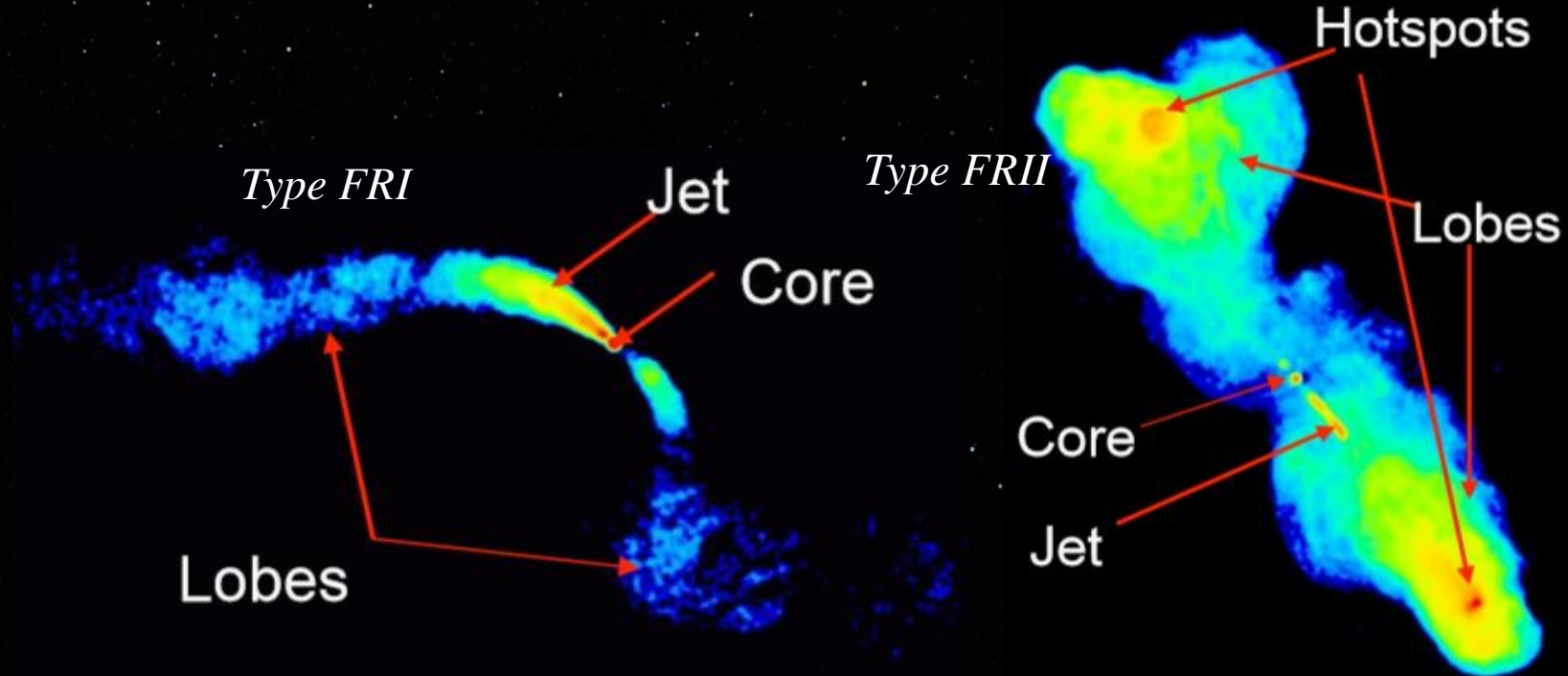
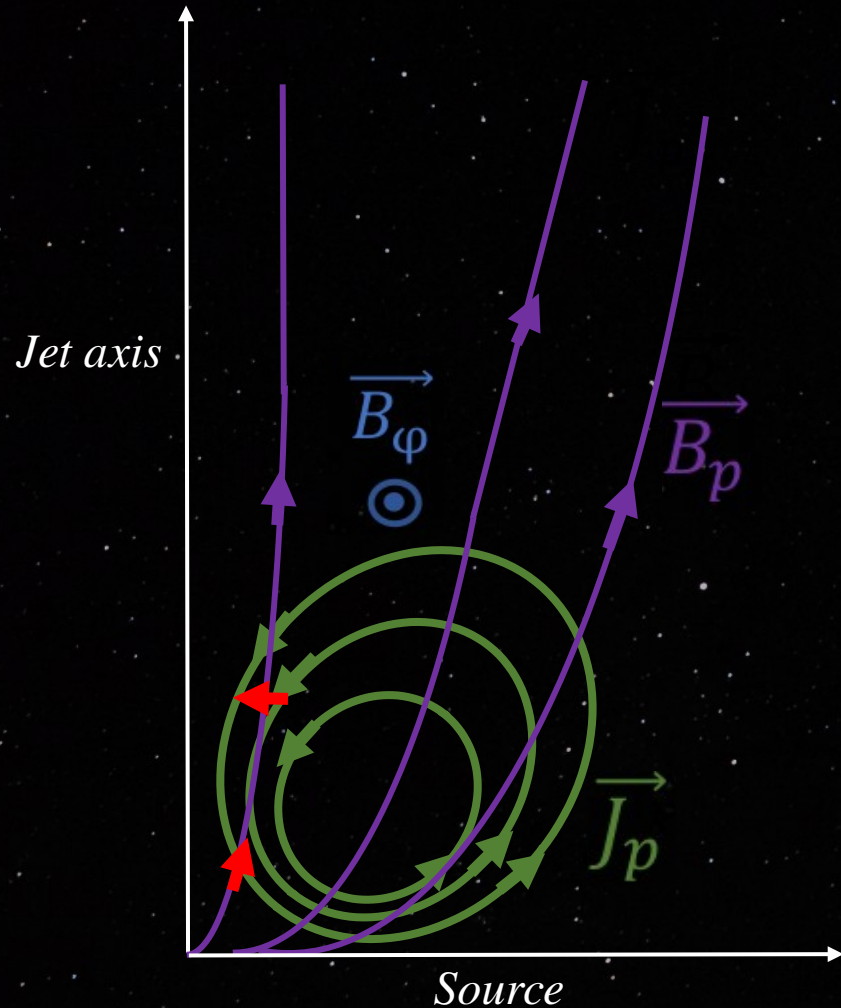
Collimation properties

- Magnetic field and electric current \Rightarrow Acceleration and collimation



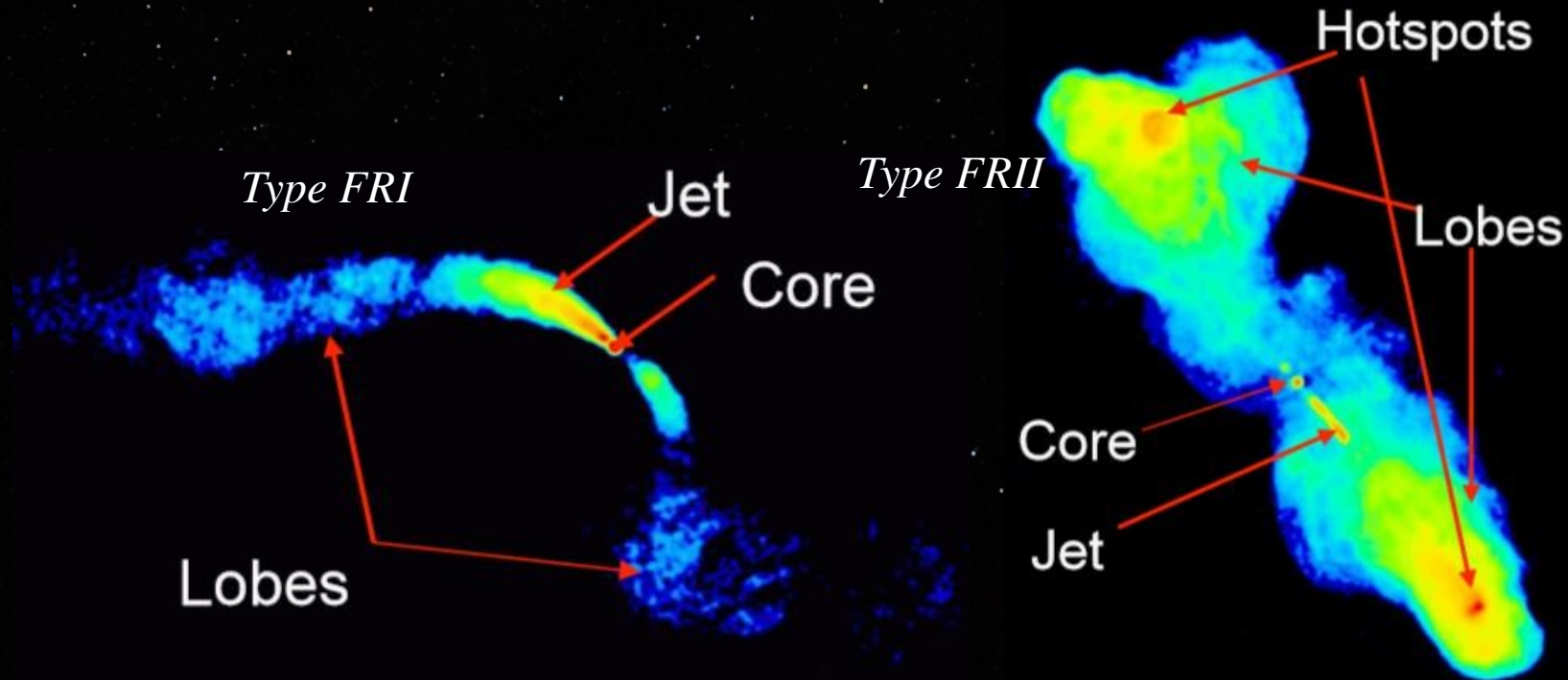
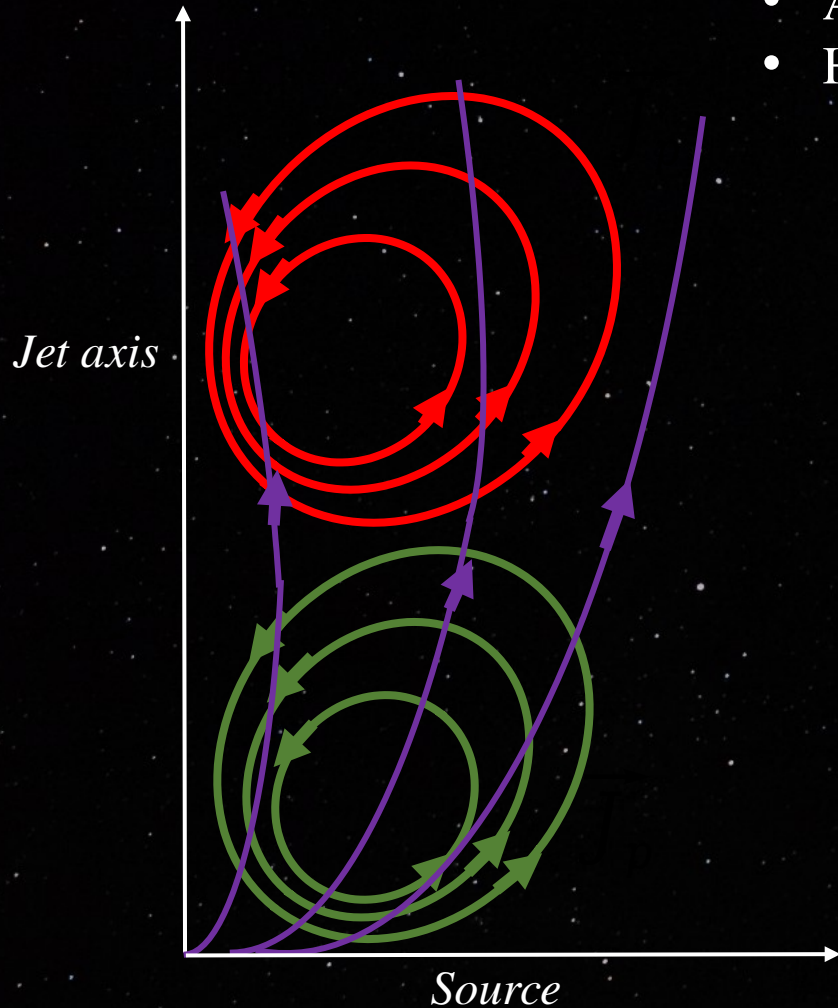
Collimation properties

- Magnetic field and electric current \Rightarrow Acceleration and collimation



Collimation properties

- Magnetic field and electric current \Rightarrow Acceleration and collimation
- Asymptotic current \Rightarrow Observable shape (Heyvaerts & Norman 1989; 2003)
- How is the asymptotic current linked to the source ?



Jet source

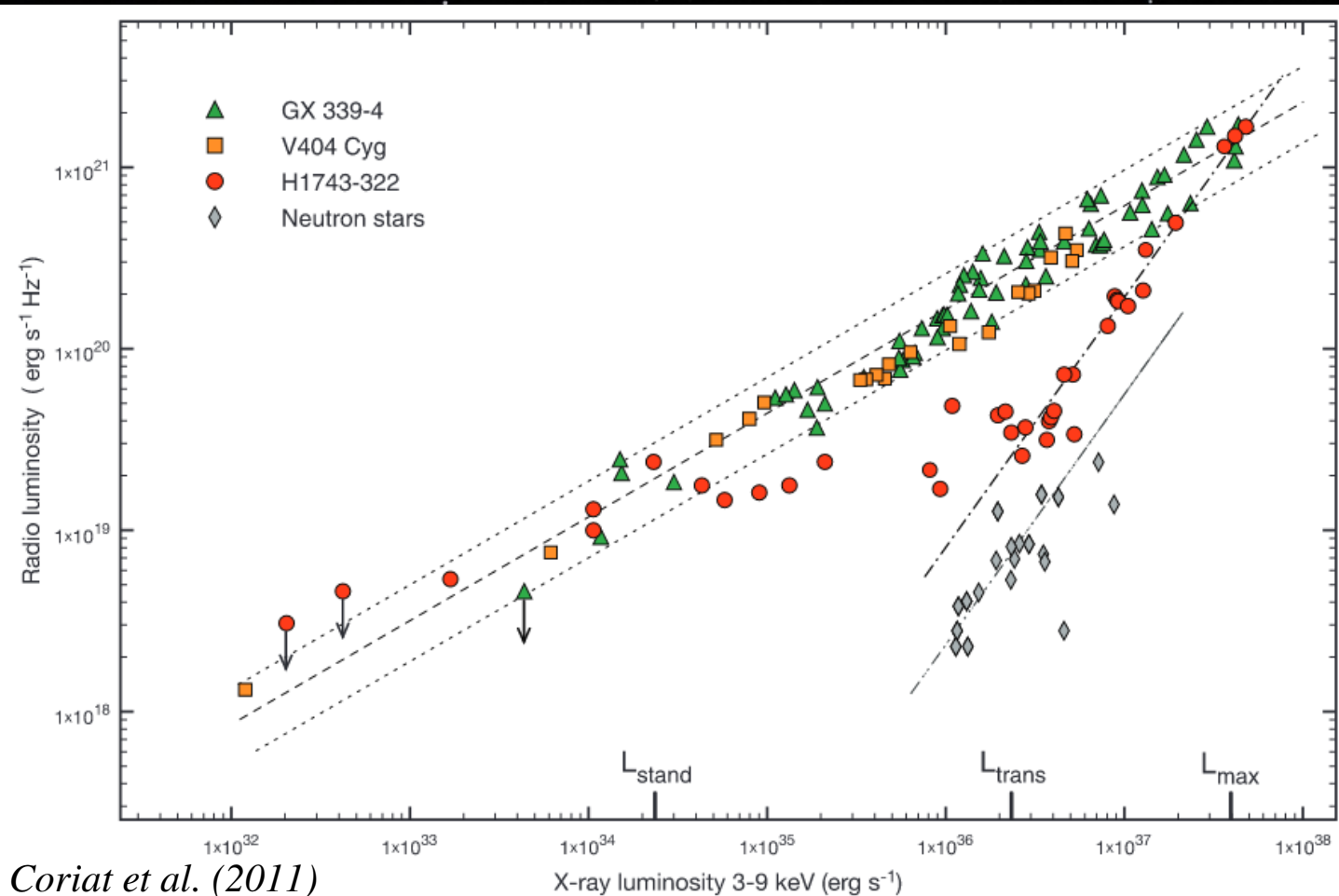
What kind of jet source ?

- Black hole (Blandford & Znajek 1977)
- Accretion disk (Blandford & Payne 1982) \Rightarrow Jet-Emitting Disk (JED, Ferreira 1997)

Jet source

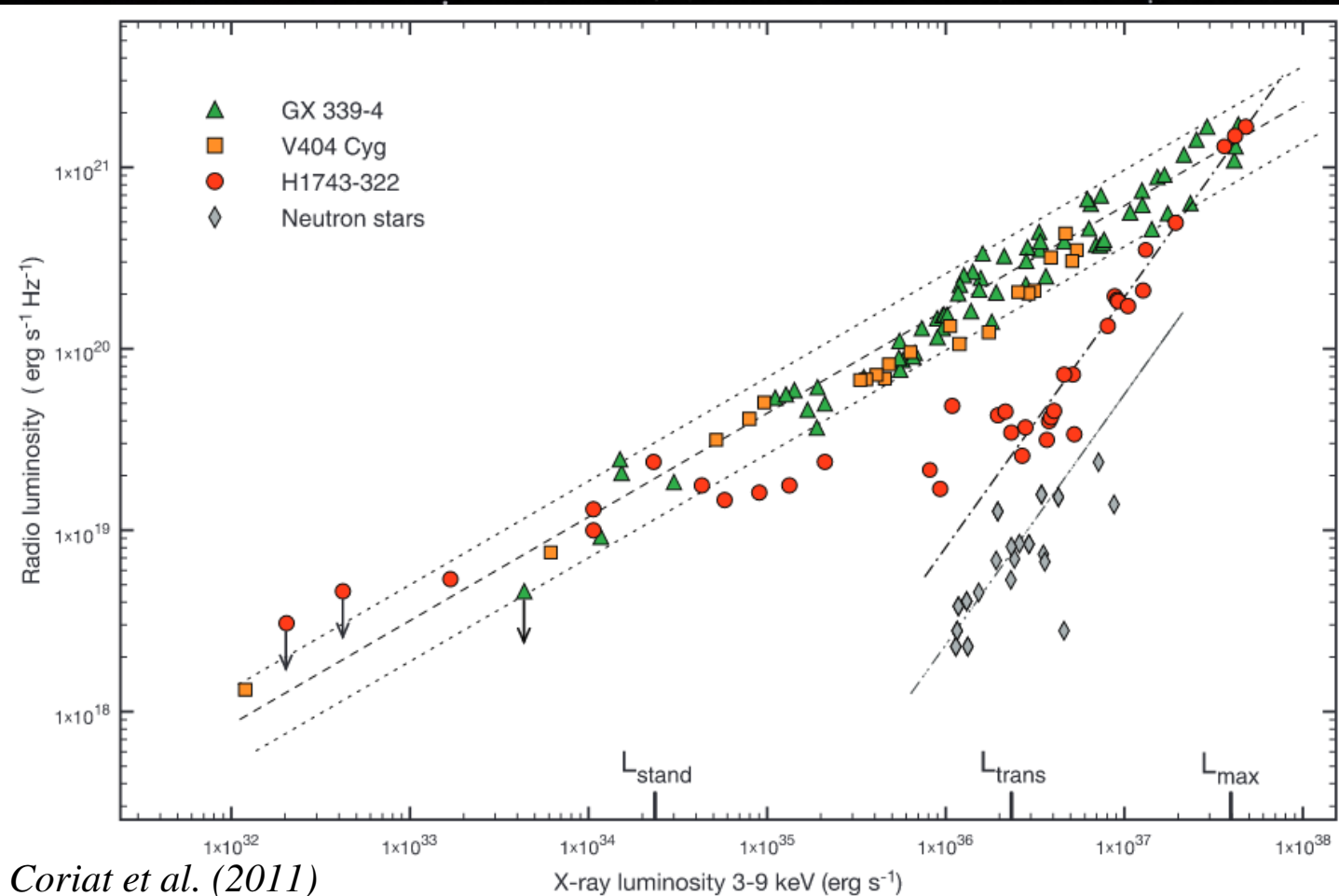
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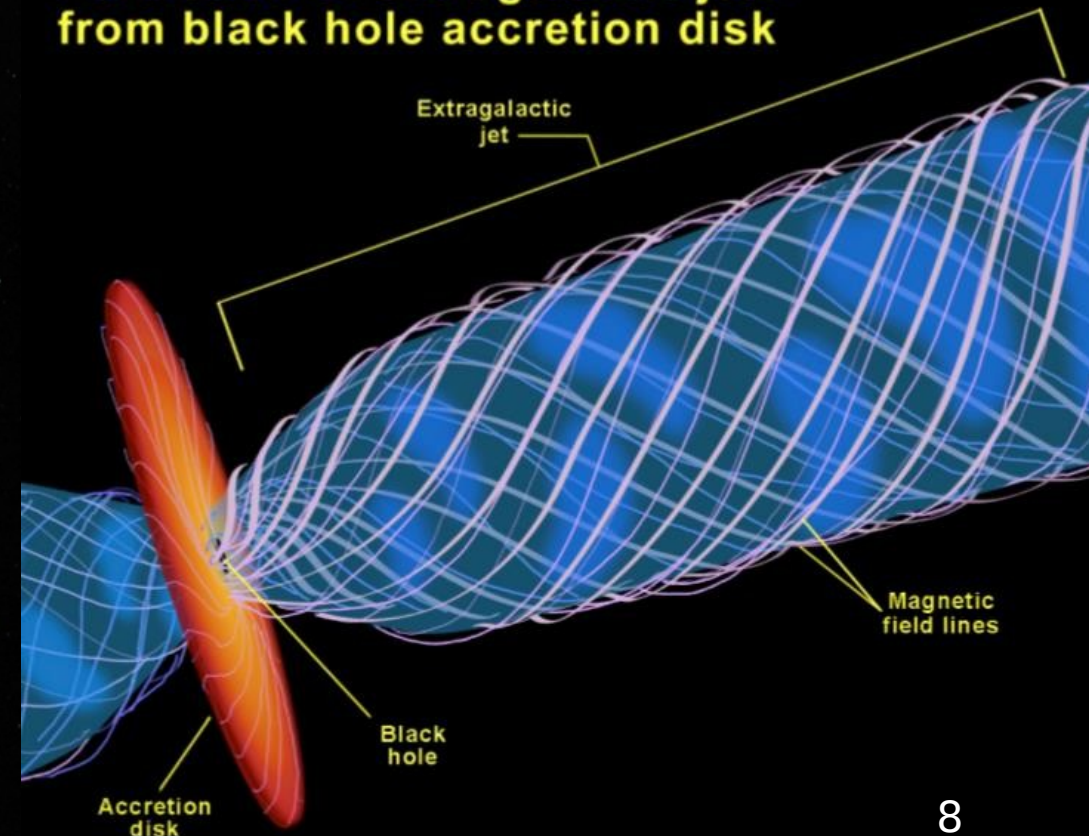
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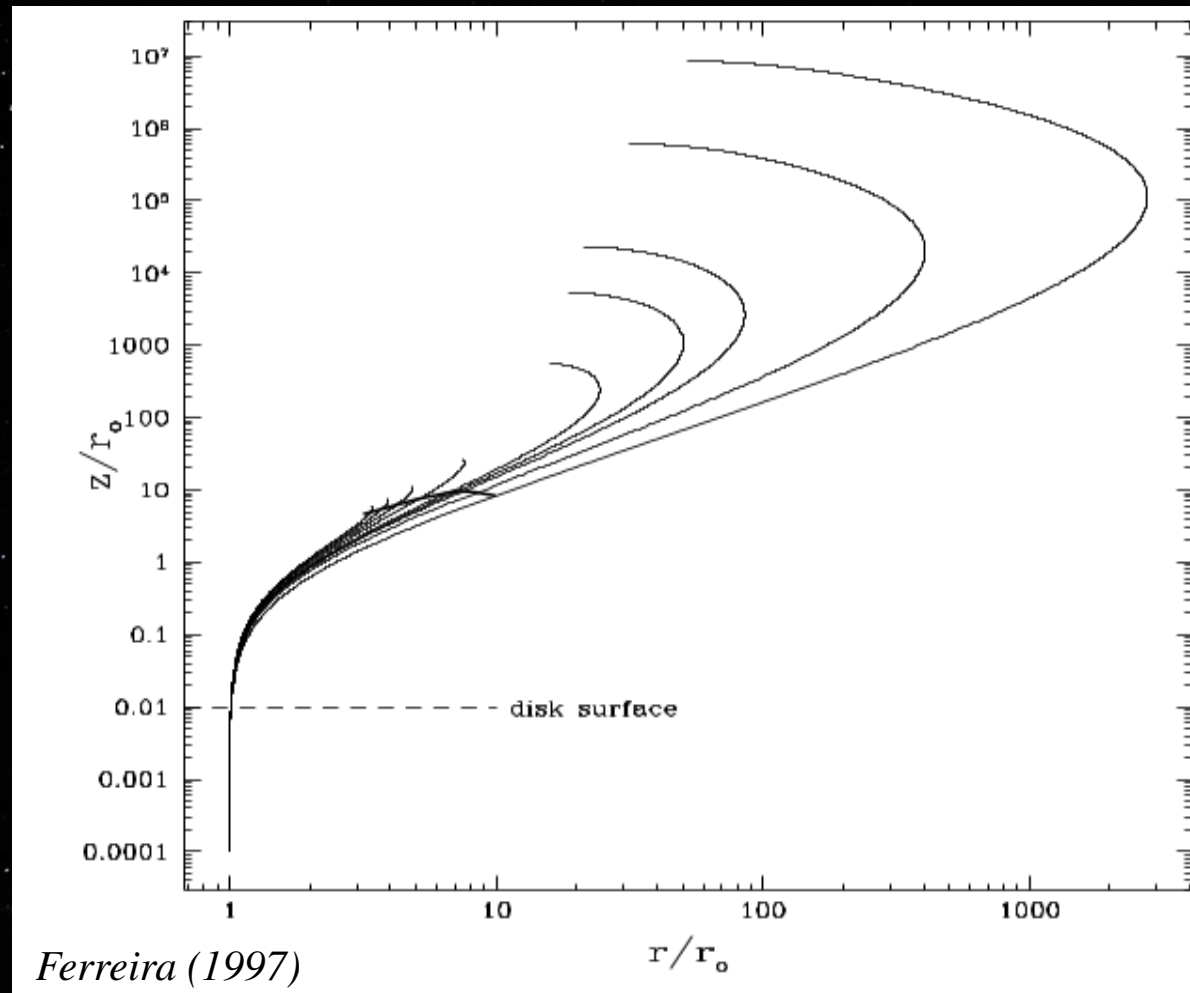
Coriat et al. (2011)

Formation of extragalactic jets from black hole accretion disk



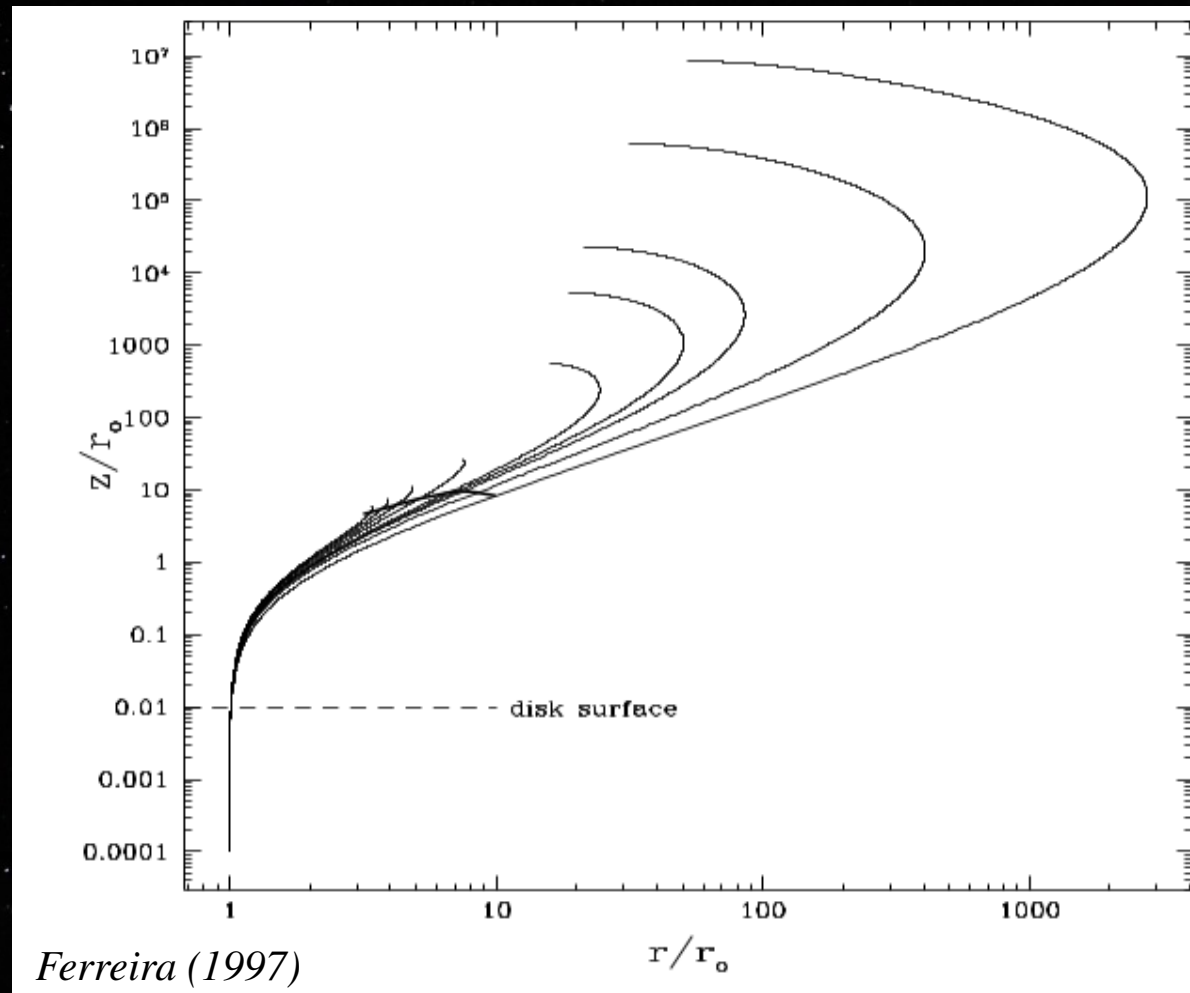
Collimation of jets emitted from JEDs

- Self-similar ansatz (Blandford & Payne 1982, Contopoulos & Lovelace 1994, Ferreira 1997)



Collimation of jets emitted from JEDs

- Self-similar ansatz (Blandford & Payne 1982, Contopoulos & Lovelace 1994, Ferreira 1997)
- Recollimation present \Rightarrow shocks expected (Ferreira 1997, Polko & al. 2010)

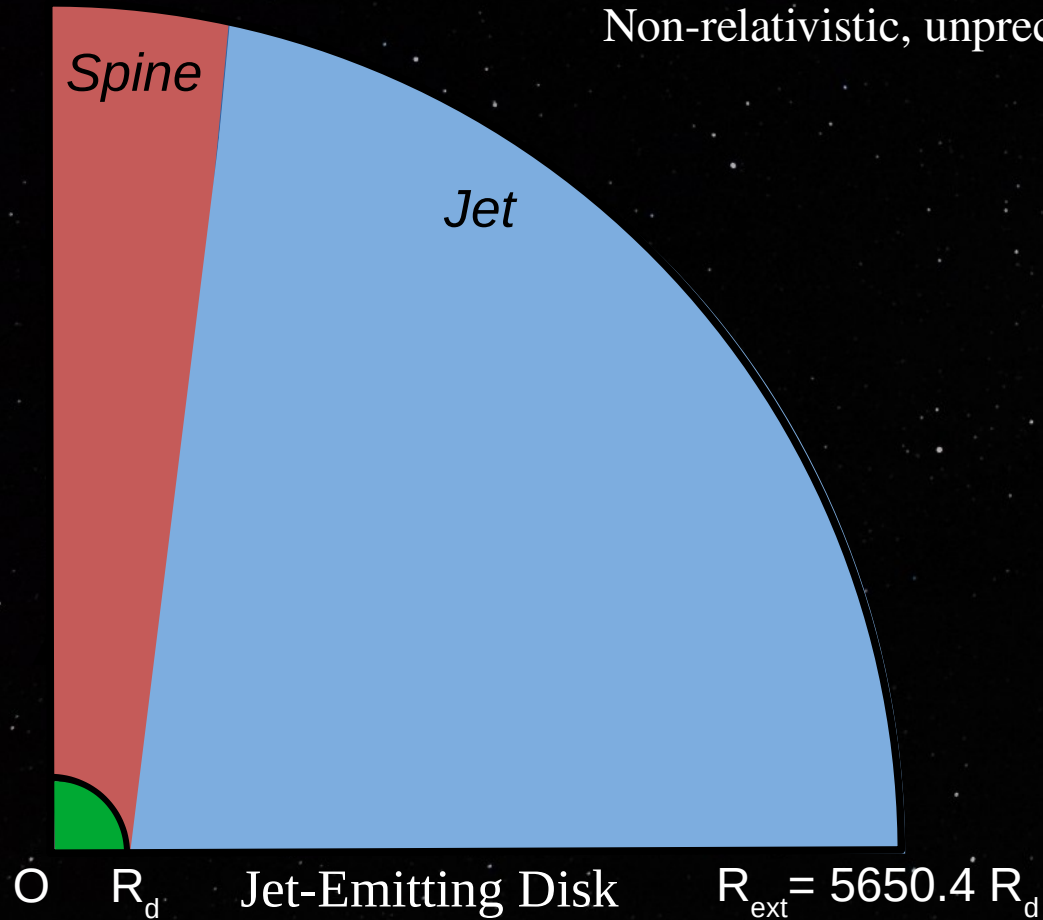


Numerical MHD simulations

Two-component outflow :

- Inner spine : Non-rotating central object (Schwarzschild Black Hole)
- Jet : Emission as in Ferreira 1997 (see Marcel+ 2018ab; 2019; 2020; 2021) on X-ray binaries

Non-relativistic, unprecedented scales



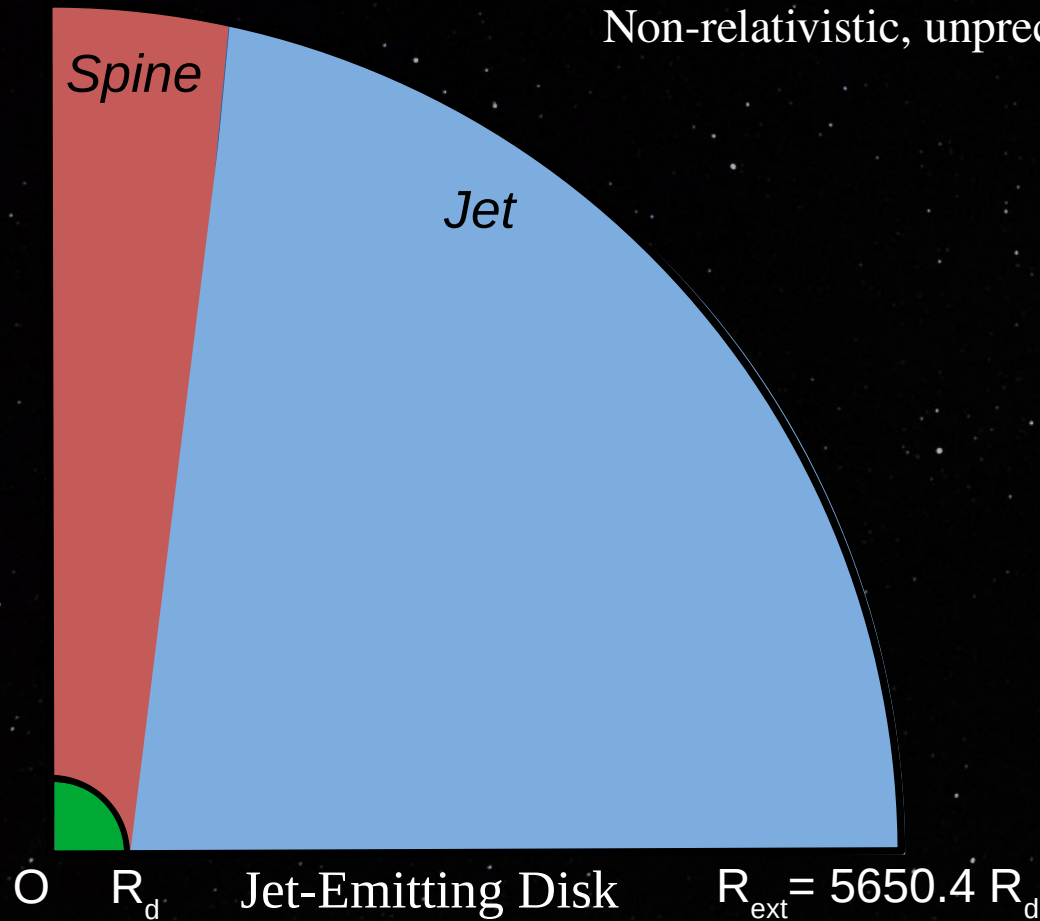
Numerical setup

Numerical MHD simulations

Two-component outflow :

- Inner spine : Non-rotating central object (Schwarzschild Black Hole)
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Non-relativistic, unprecedented scales

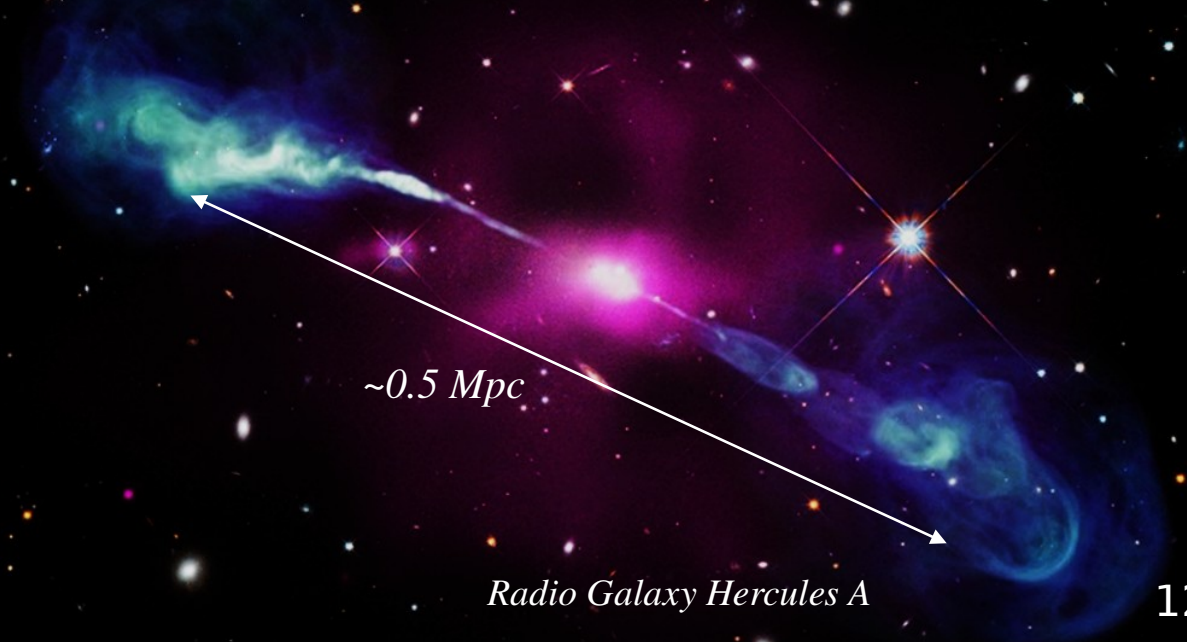


For a Schwarzschild black hole, $R_d = R_{\text{ISCO}} = 6 \frac{GM}{c^2}$

For a $10^{10} M_{\odot}$ black hole :

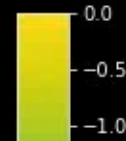
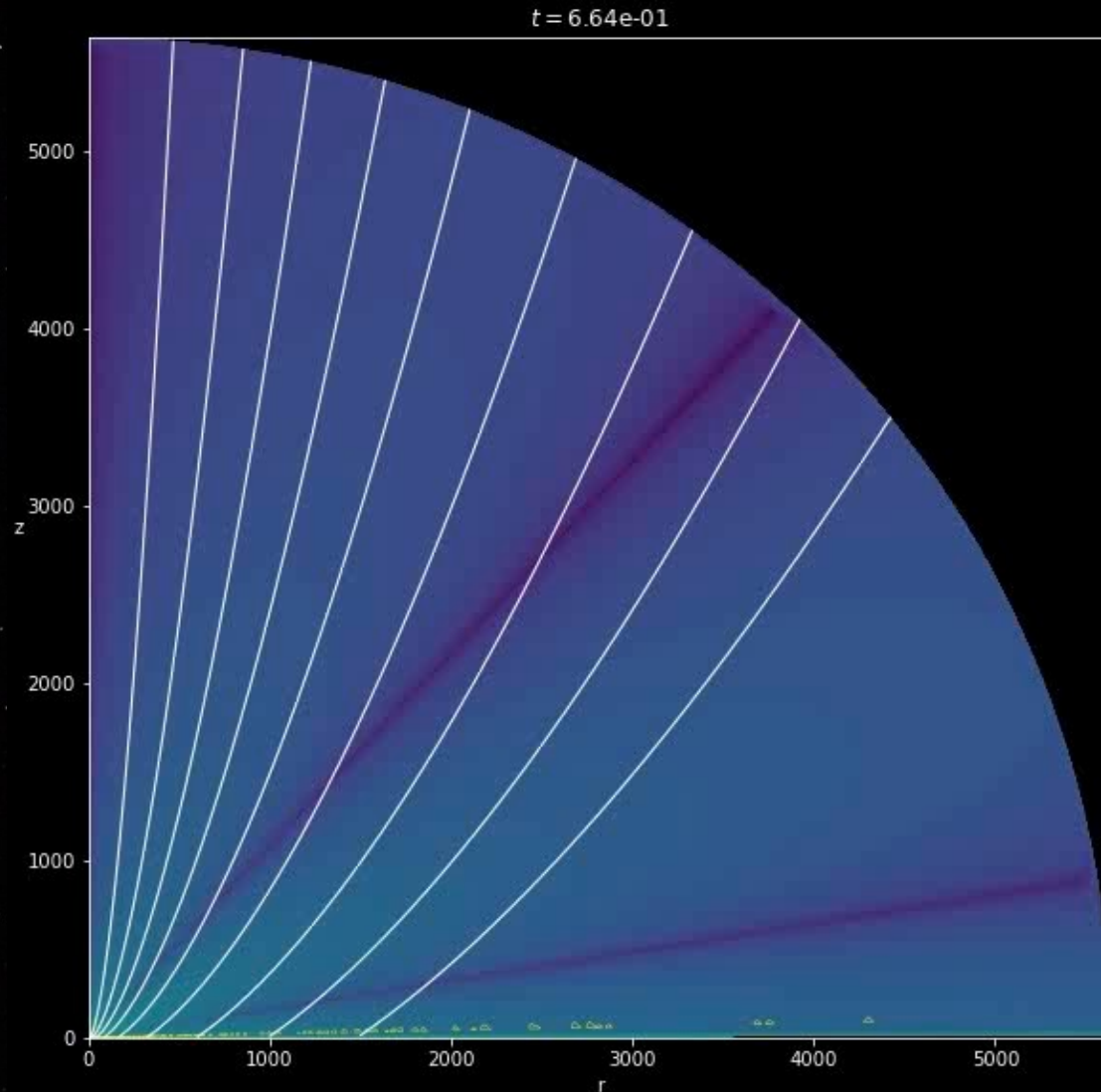
$$R_d = 5 \cdot 10^{-3} \text{ pc}$$

$$R_{\text{ext}} = 30 \text{ pc}$$



Radio Galaxy Hercules A

Numerical MHD simulations



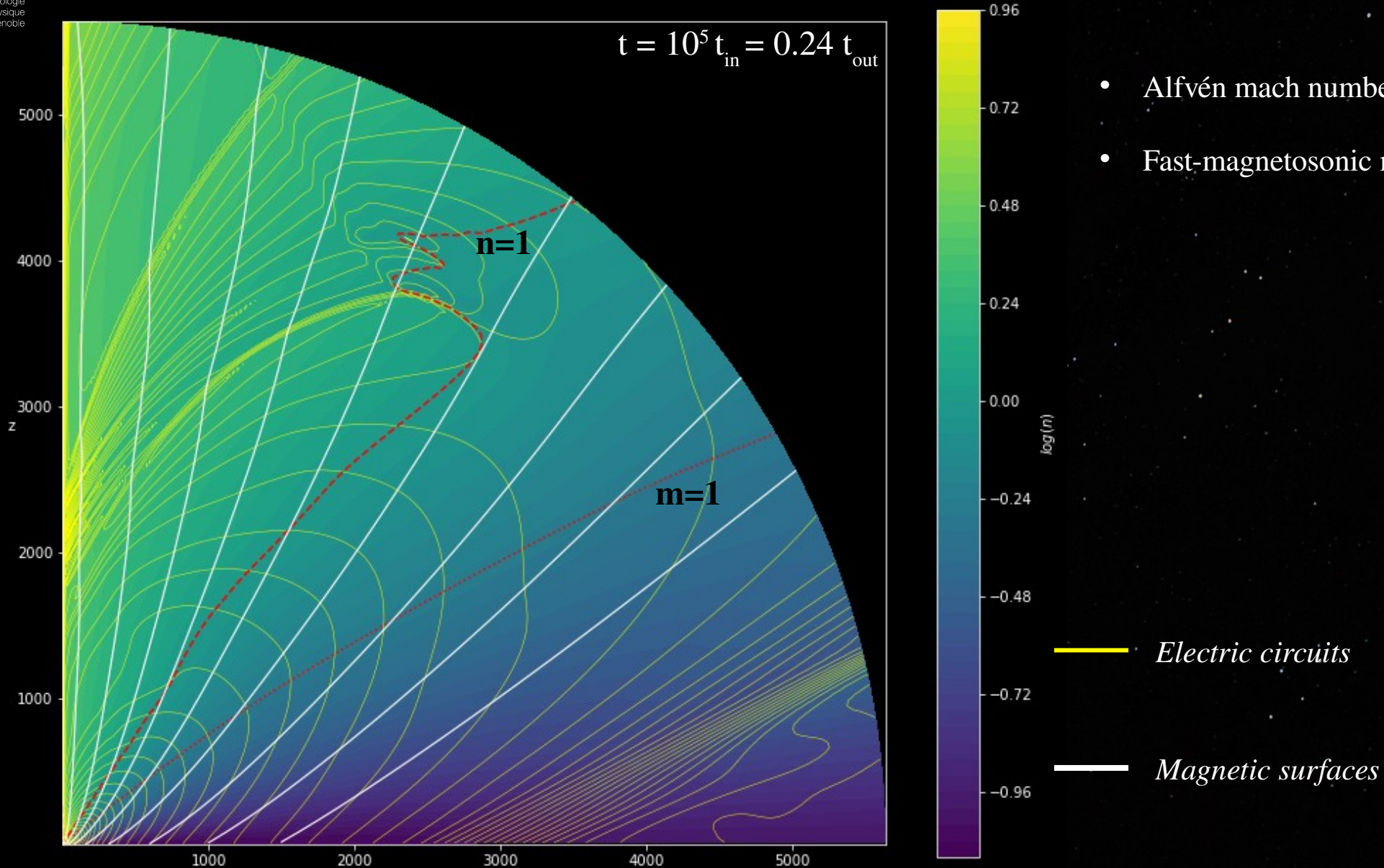
$\log(n)$

- Alfvén mach number : $m = \frac{u_p}{v_{A_p}}$
- Fast-magnetosonic mach number : $n = \frac{u_p}{v_{FM_p}}$

 *Electric circuits*

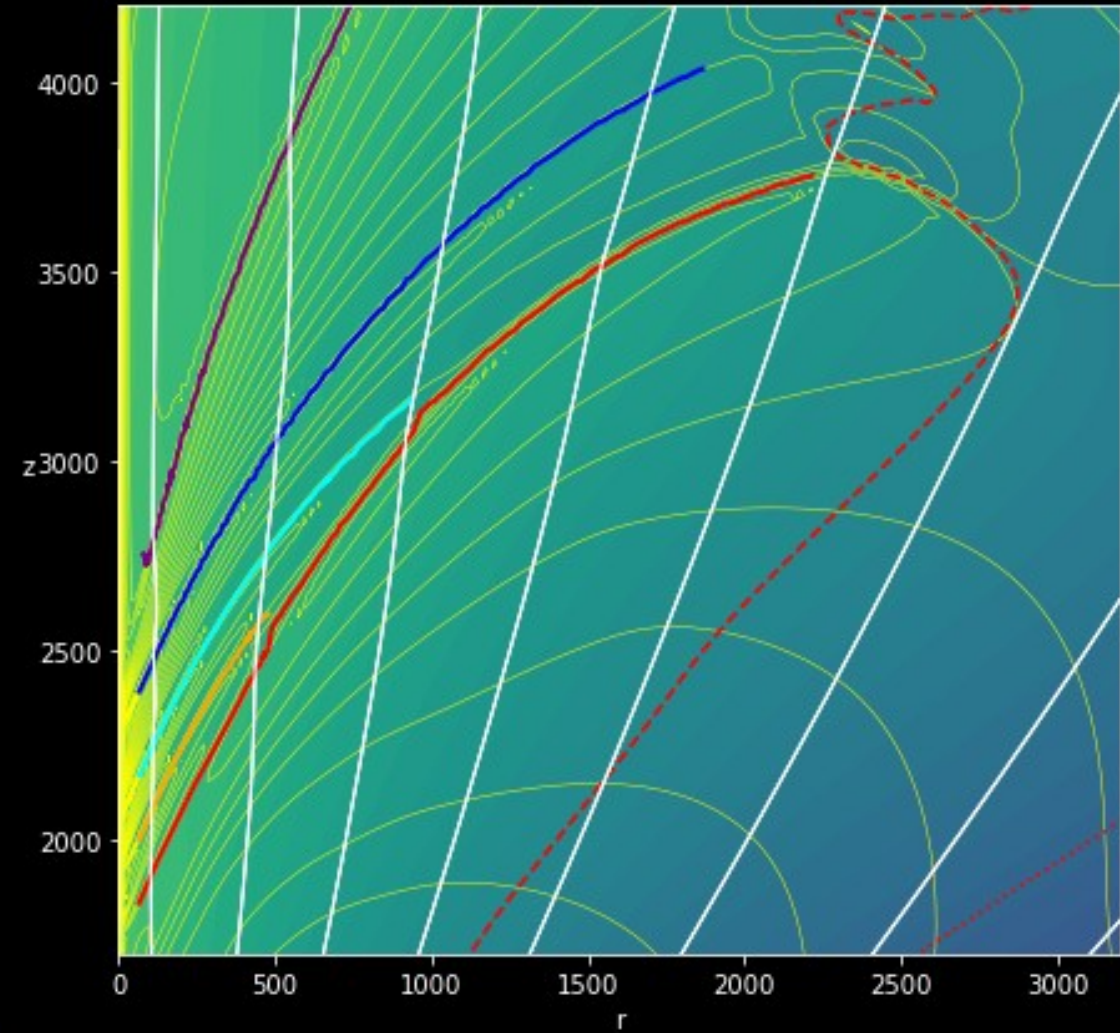
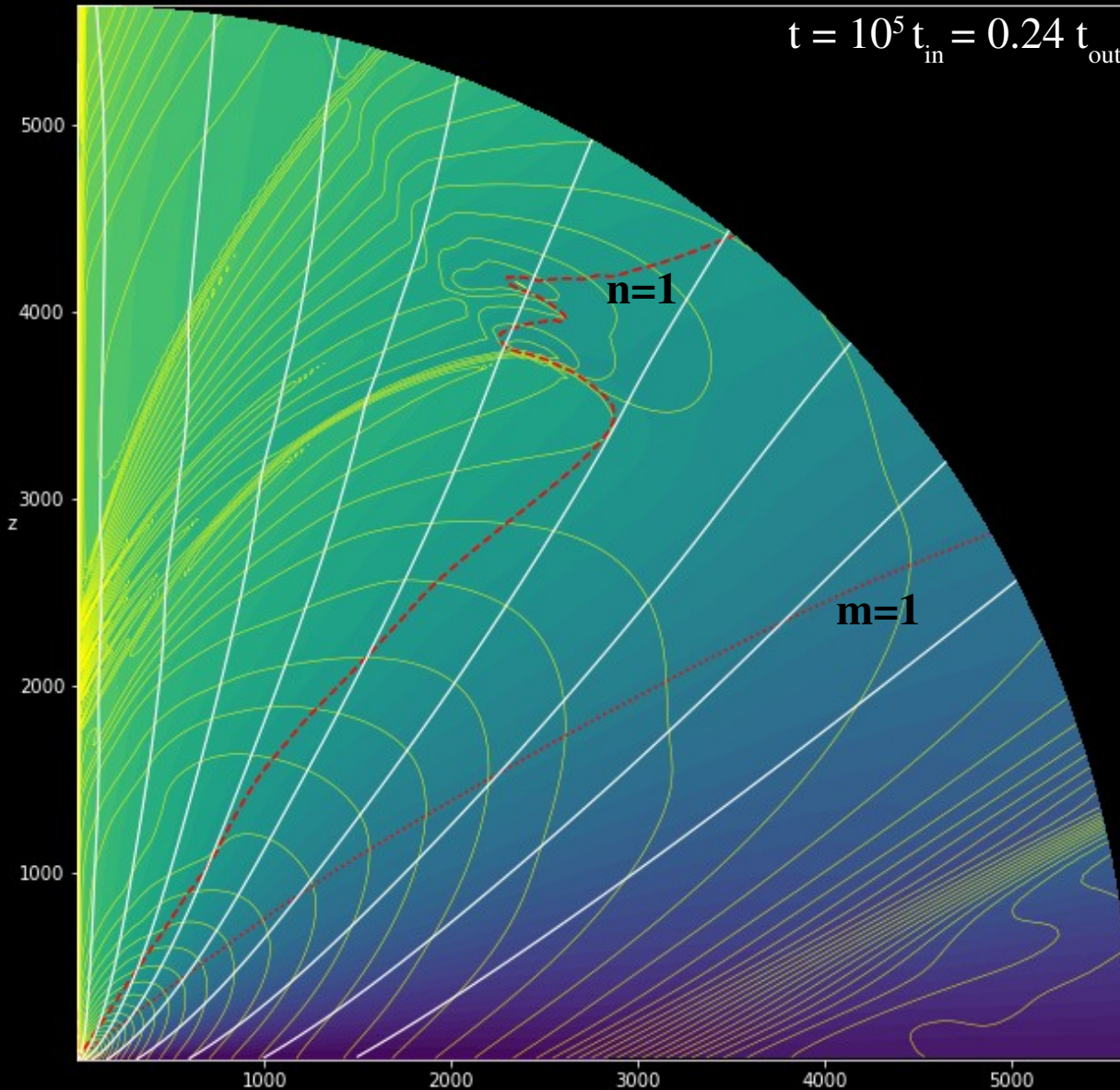
 *Magnetic surfaces*

MHD recollimation shocks



- Alfvén mach number : $m = \frac{u_p}{v_{A_p}}$
- Fast-magnetosonic mach number : $n = \frac{u_p}{v_{FM_p}}$

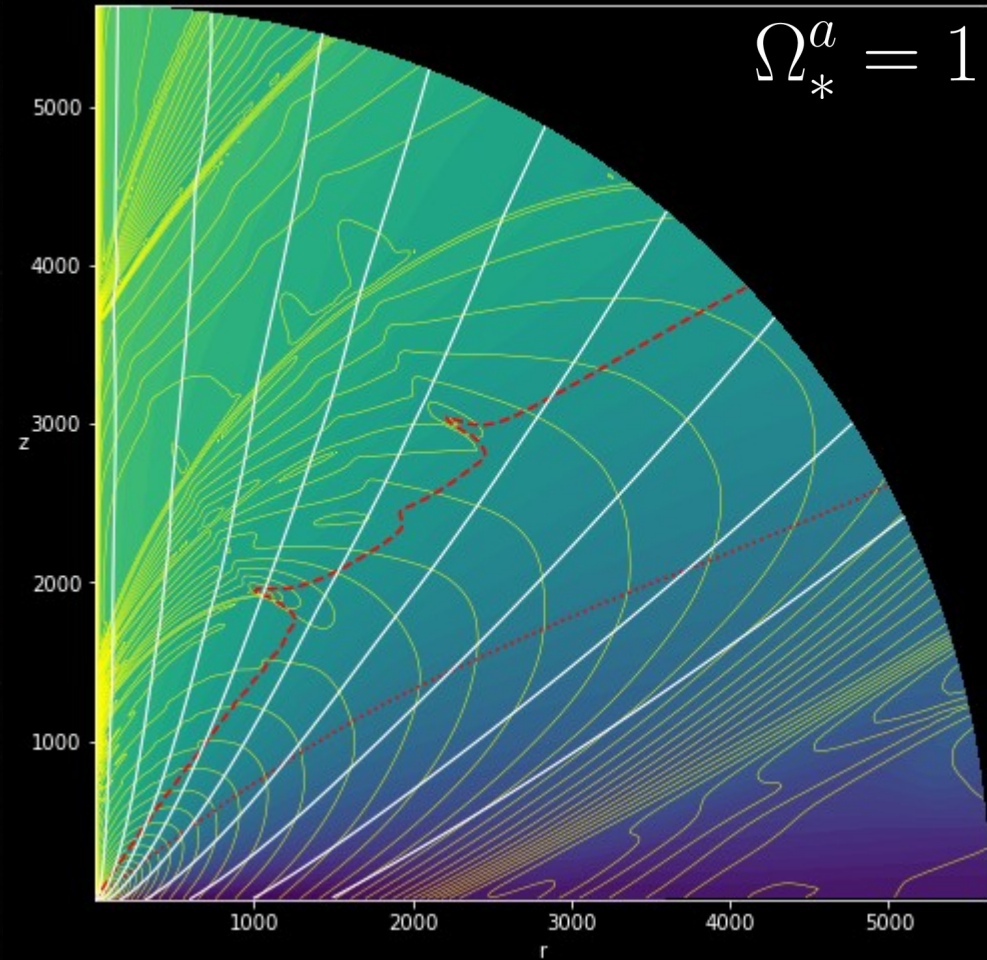
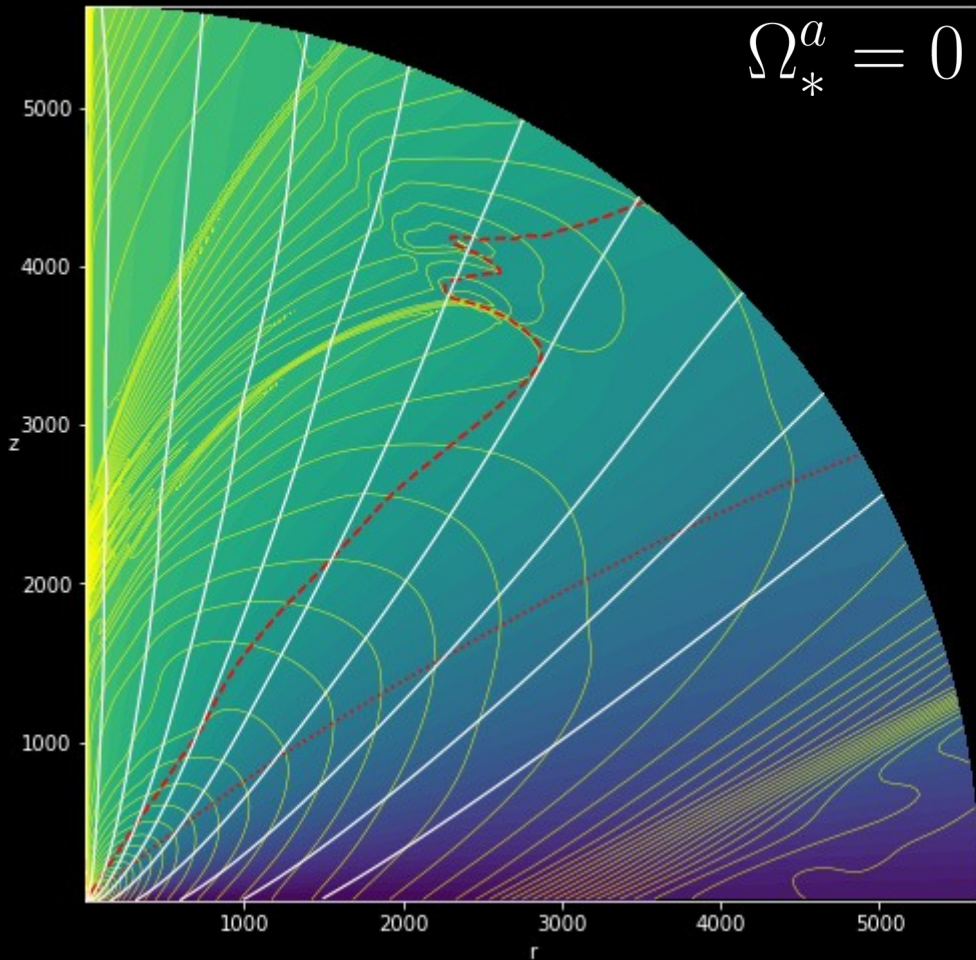
MHD recollimation shocks



Zoom on the recollimation shocks

Influence of the spine

- Apathic simulations with the JED parameter space
- Strong influence of the central object rotation



Conclusion

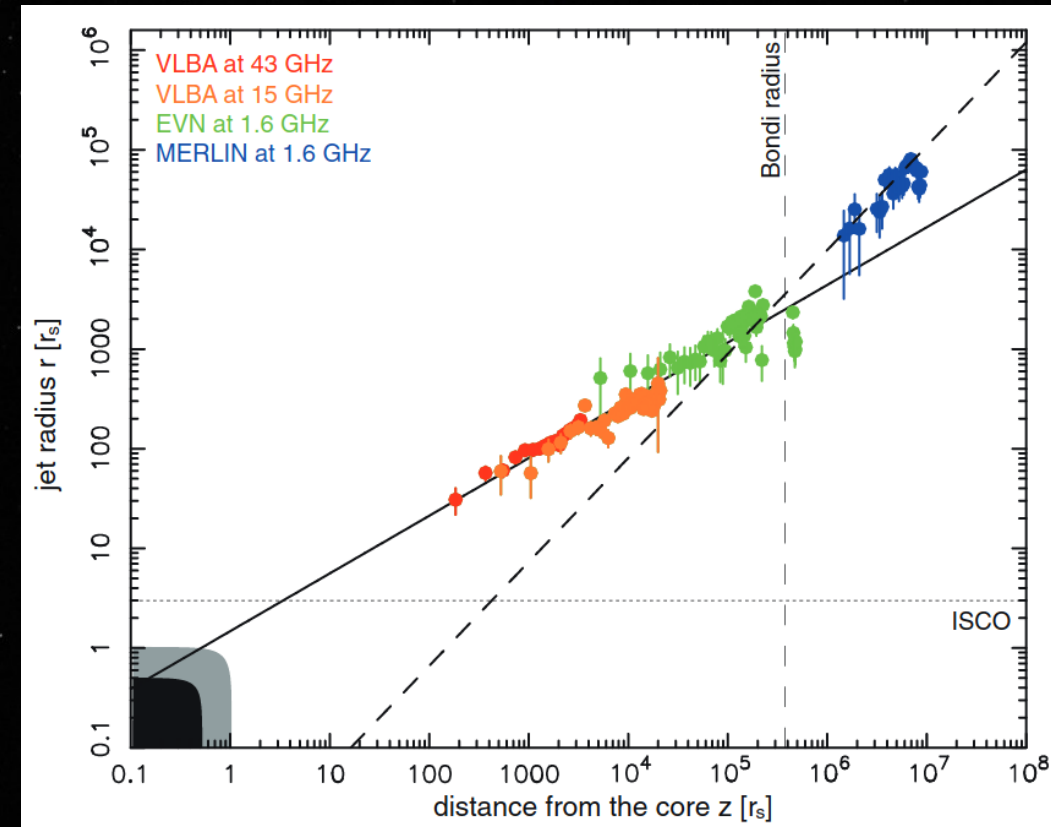
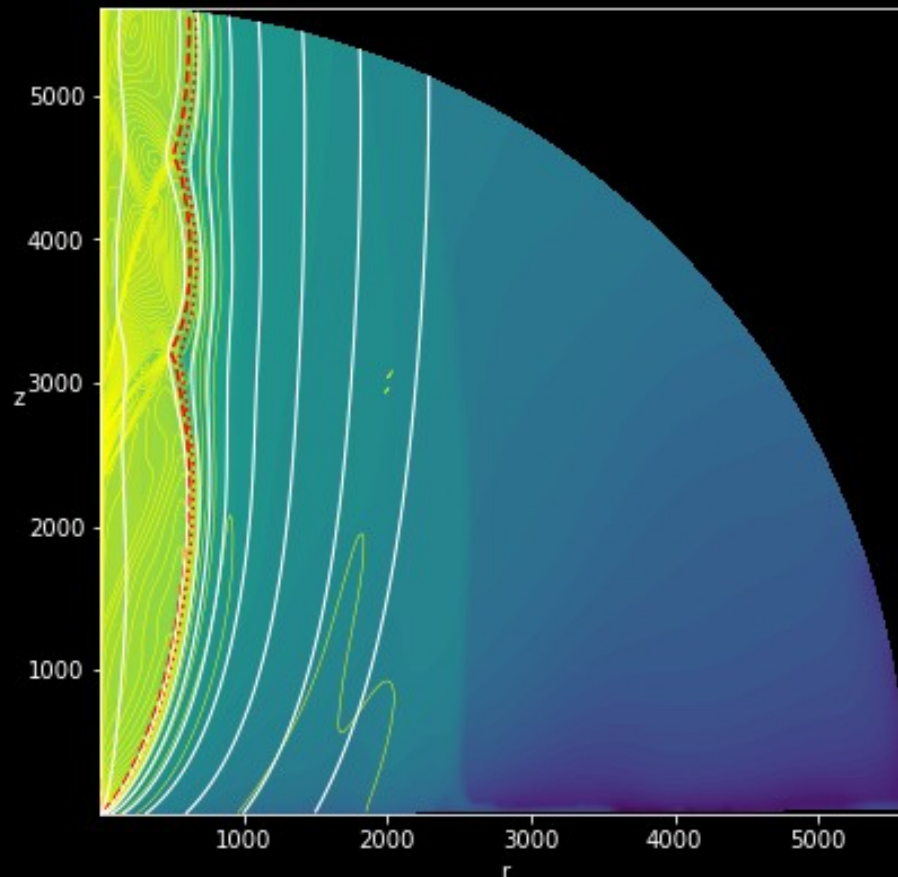
Jets emitted from JEDs of a large radial extent \Rightarrow Recollimation shocks

- 2D MHD simulations on unprecedented scales in space and time
- Jannaud, Zanni & Ferreira, submitted to A&A

Conclusion

Jets emitted from JEDs of a large radial extent \Rightarrow Recollimation shocks

- 2D MHD simulations on unprecedented scales in space and time
- Jannaud, Zanni & Ferreira, submitted to A&A
- Comparisons to observations of AGN and YSO jets \Rightarrow JEDs of finite size

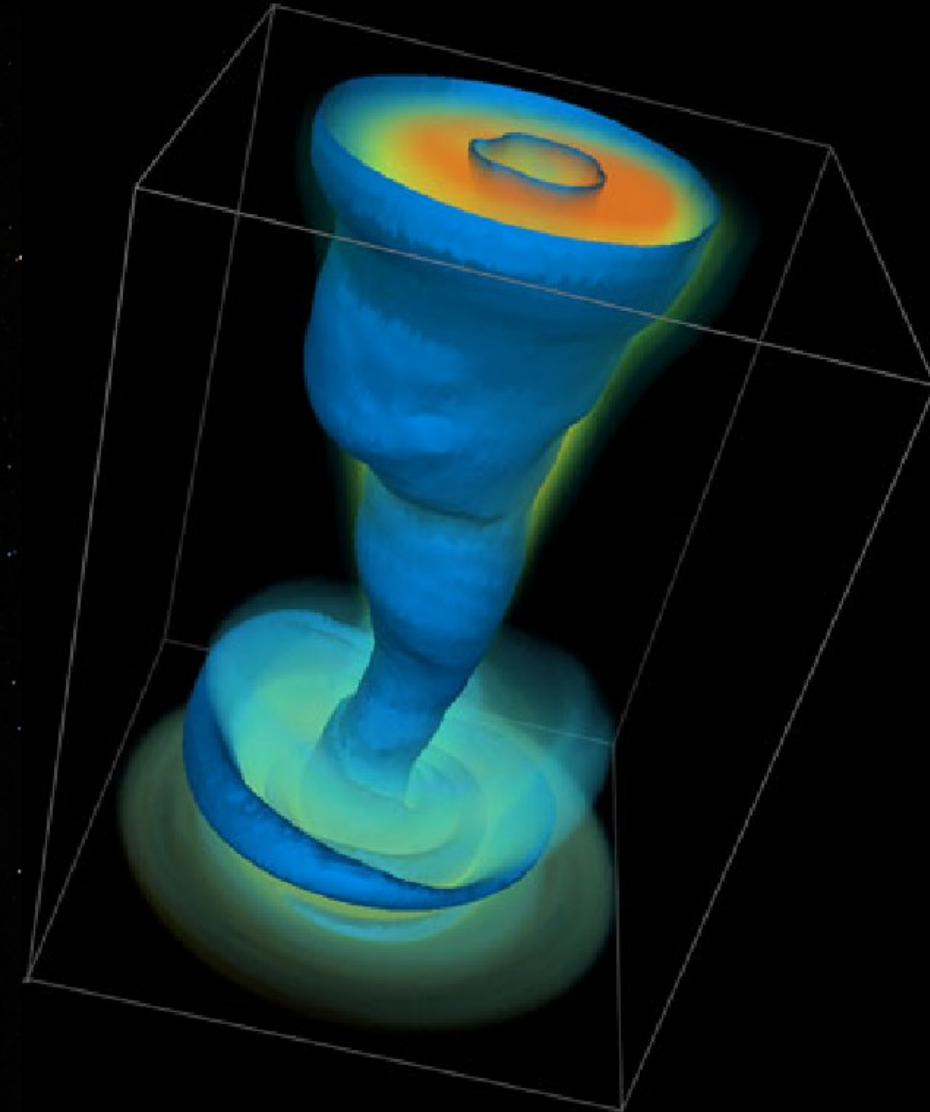
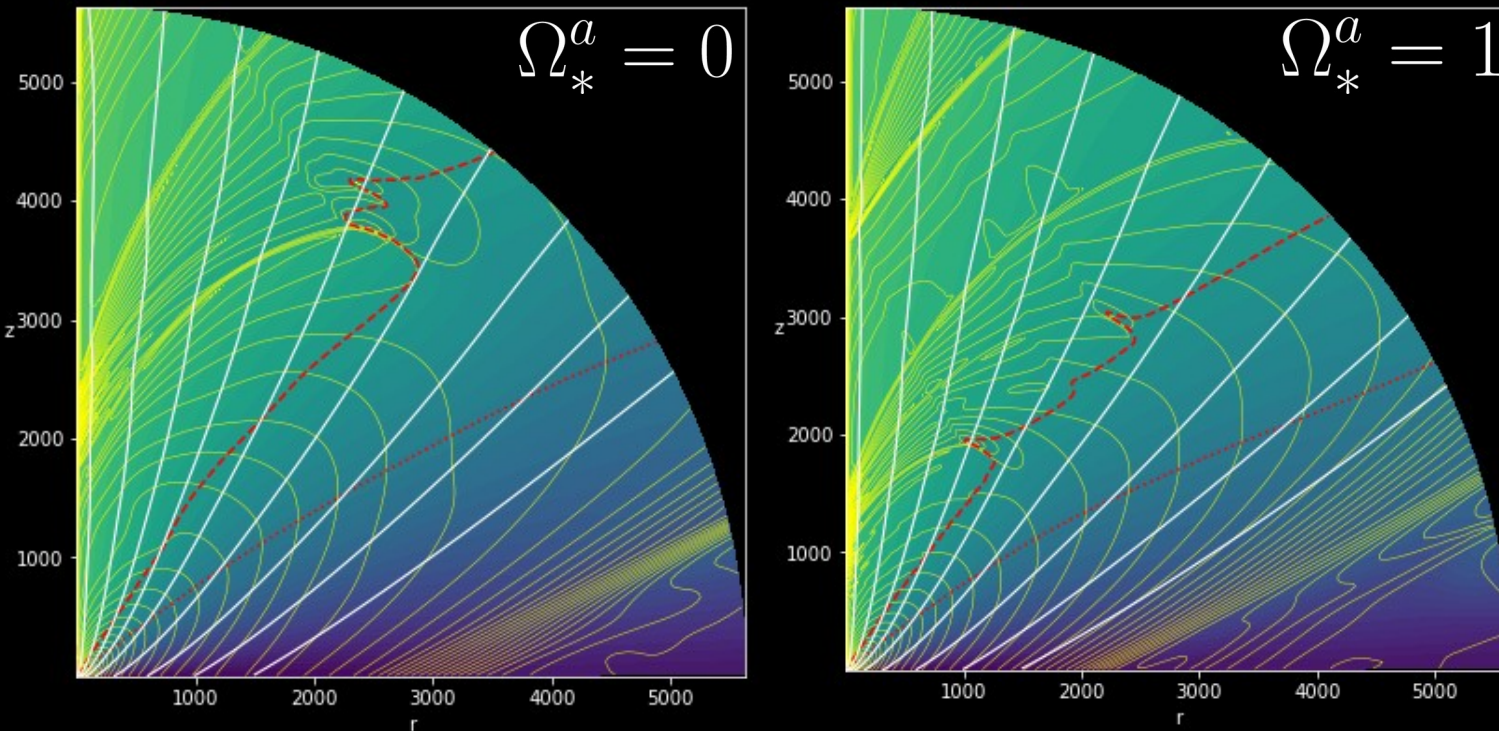


HST-1 complex in M87, Asada & Nakamura (2012)

Conclusion

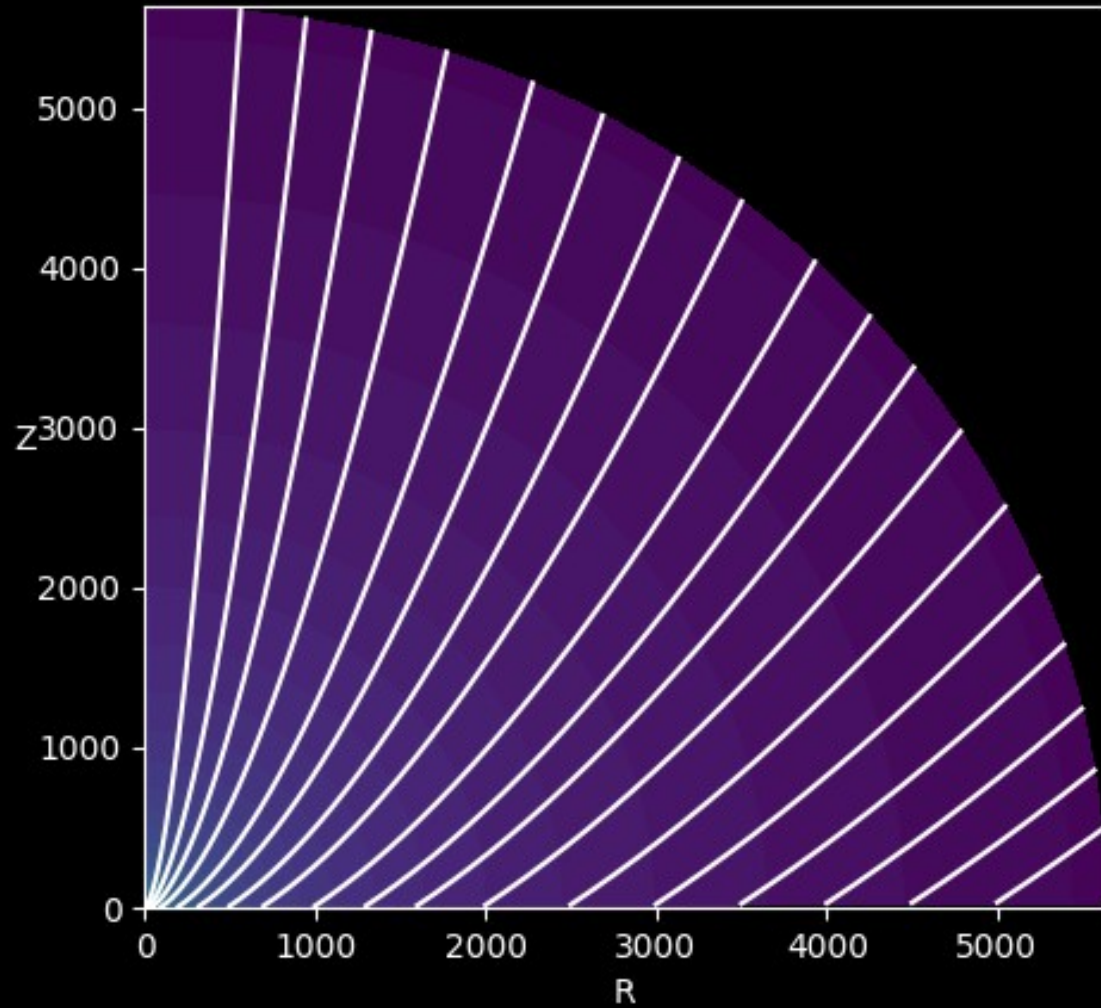
Strong influence of the spine on collimation

- Possible presence of a Blandford & Znajek outflow (see Barnier et al. 2021)
- Preliminary results (other simulations in the making, including in 3D)

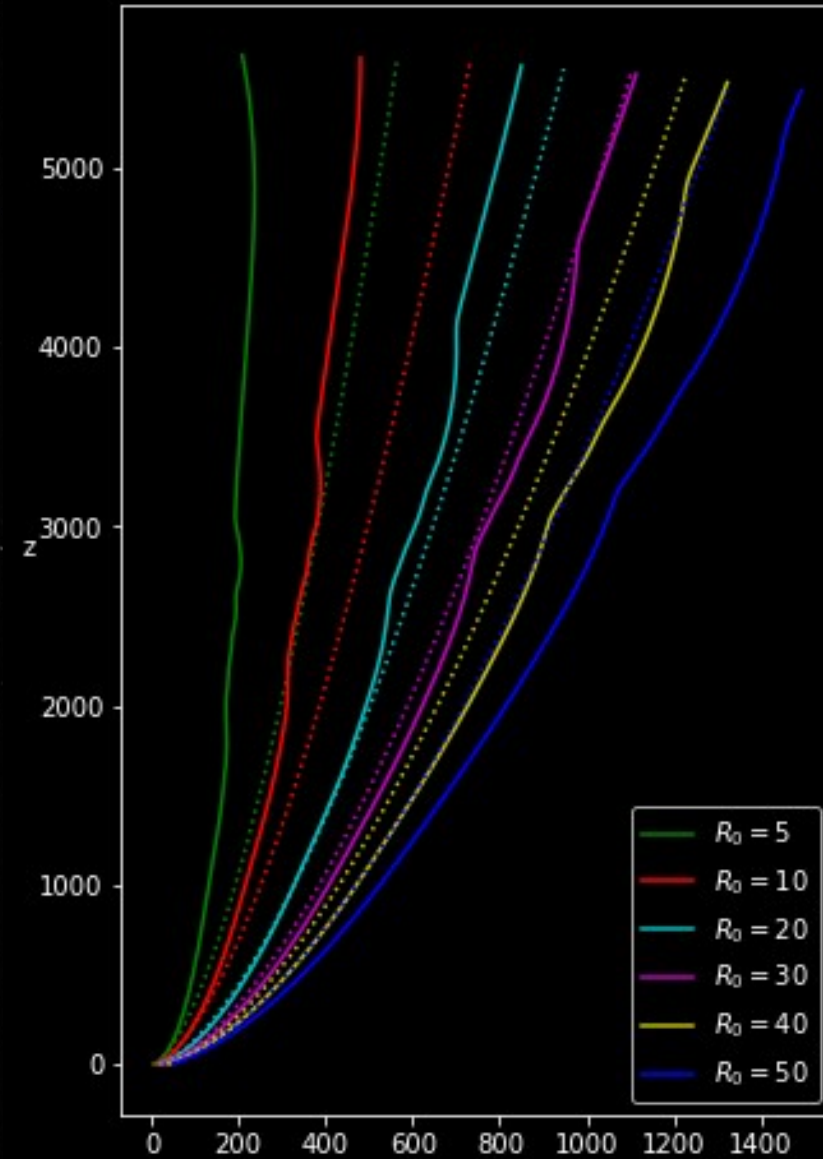


McKinney & Blandford (2009)

Collimation of the magnetic field lines

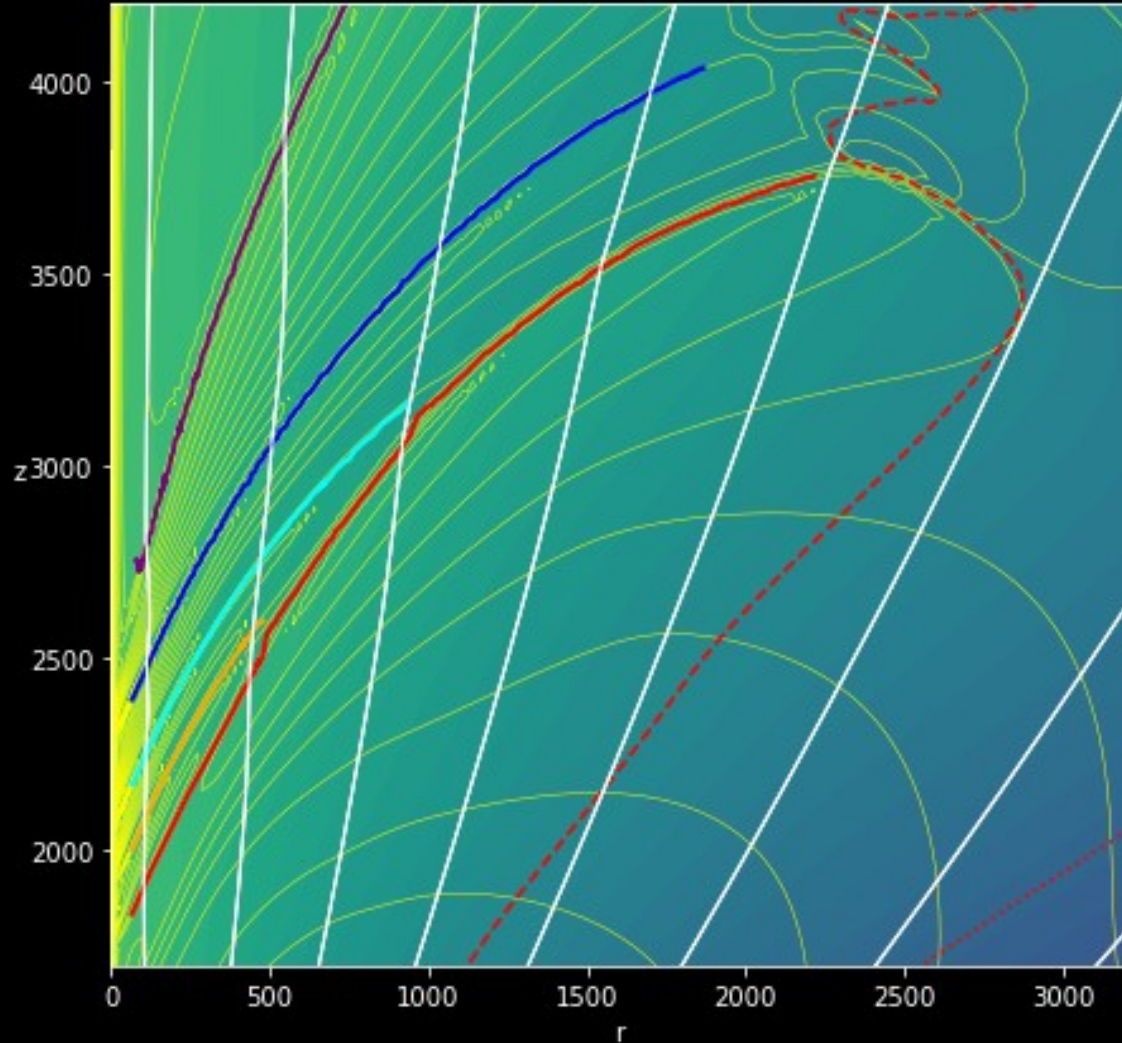


Initial conditions

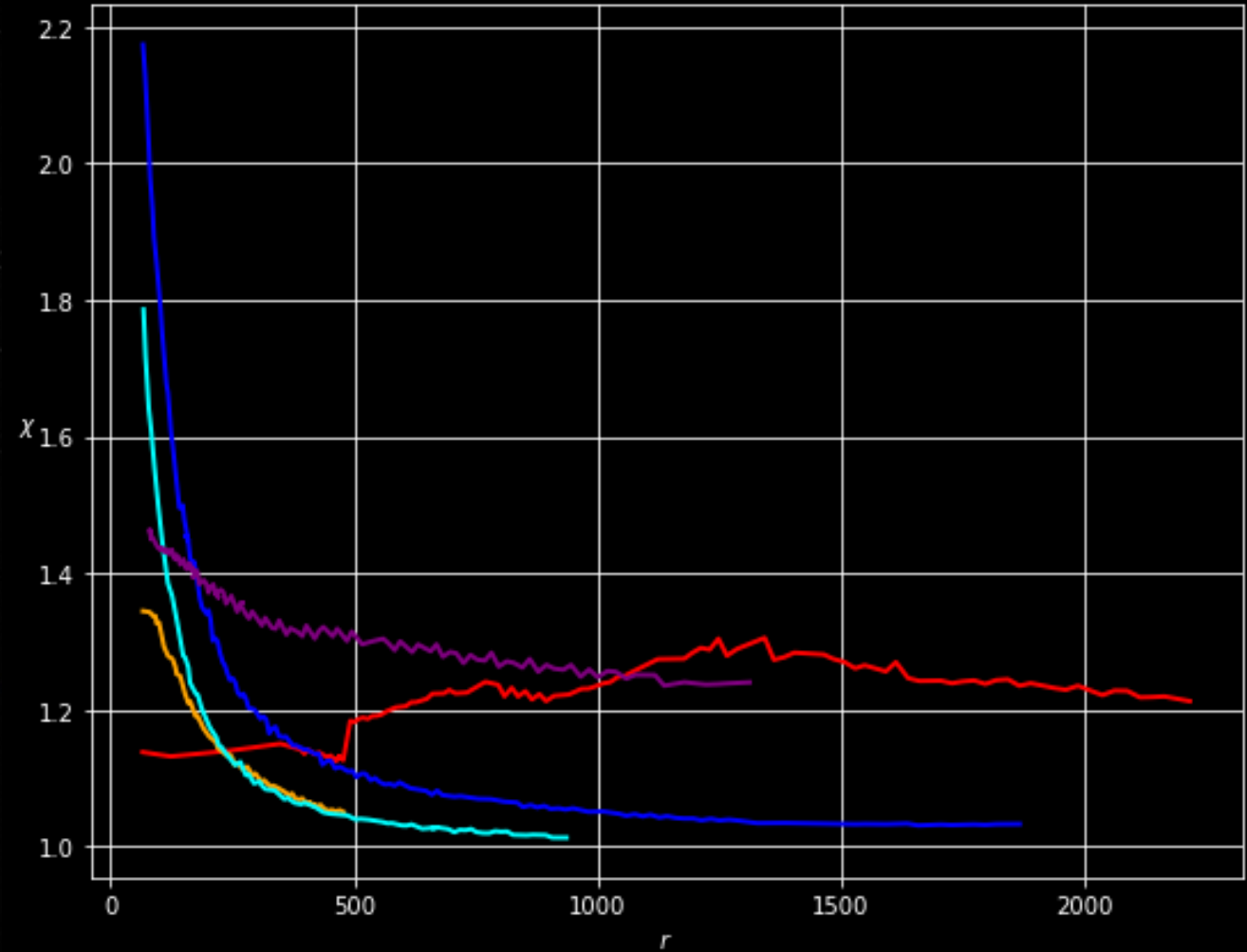


Field lines evolution

Weak and multiple recollimation shocks



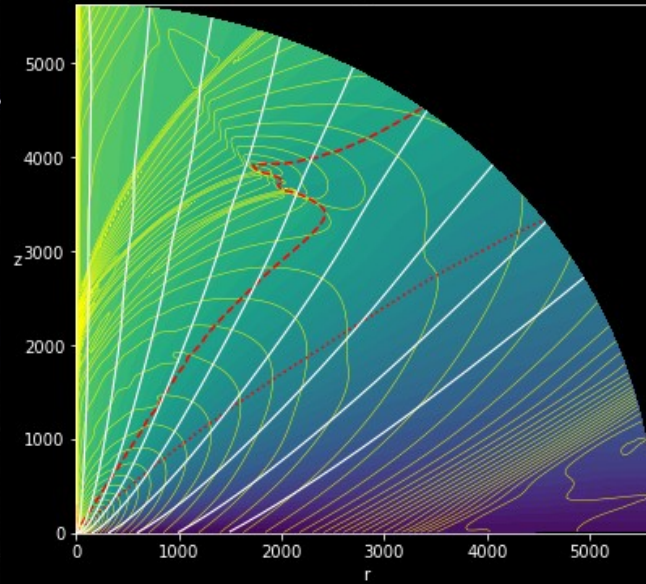
Zoom on the recollimation shocks



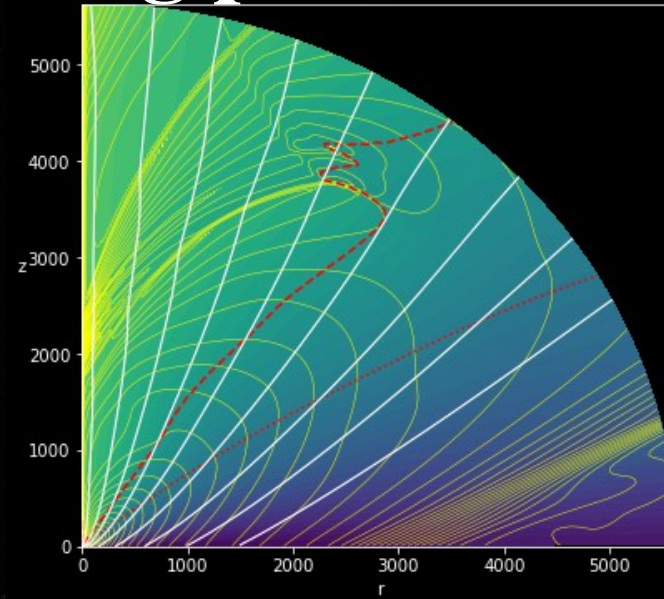
Evolution of the compression rate along the shocks

Influence of the mass loading parameter κ

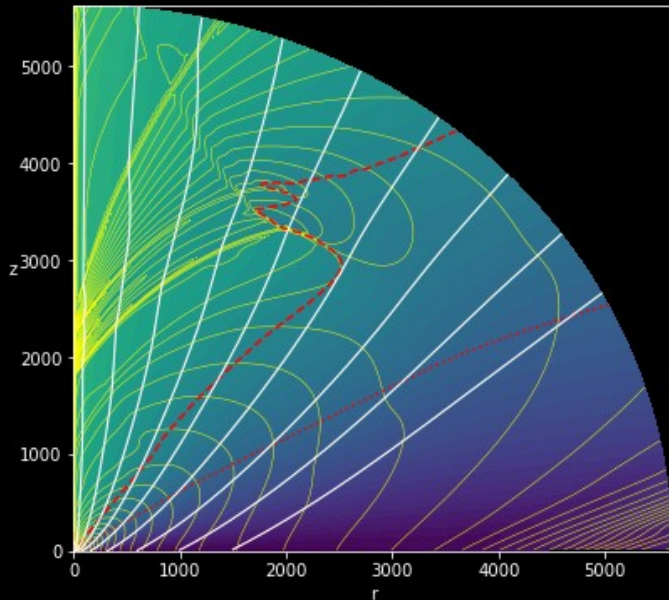
$$\kappa = \frac{v_{\theta d} v_{k d}}{v_{A d}^2}$$



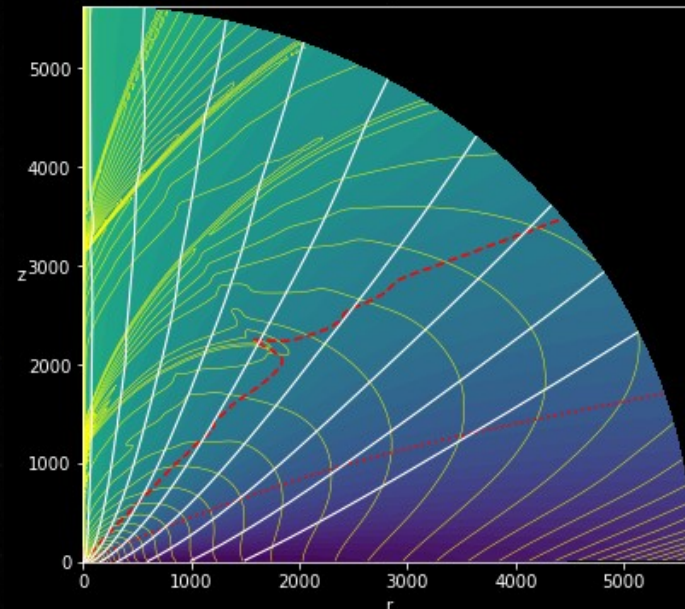
$\kappa = 0,05$



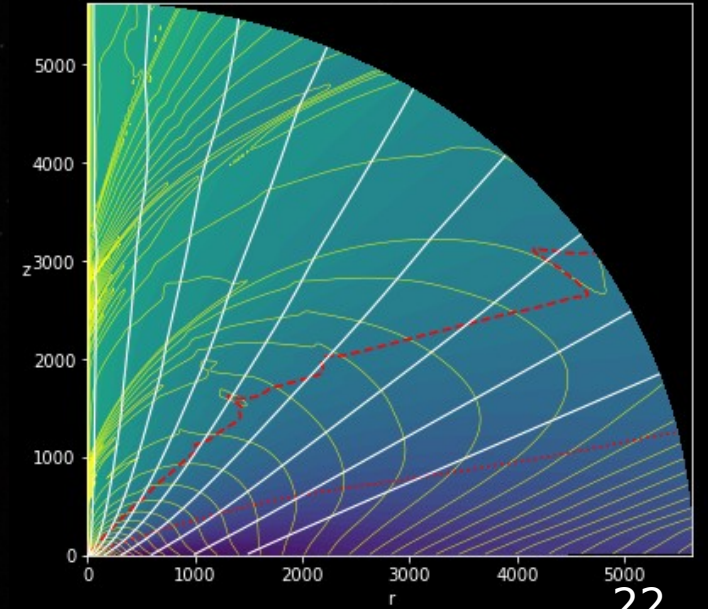
$\kappa = 0,1$



$\kappa = 0,2$



$\kappa = 0,5$



$\kappa = 1,0$

Influence of the mass loading parameter κ

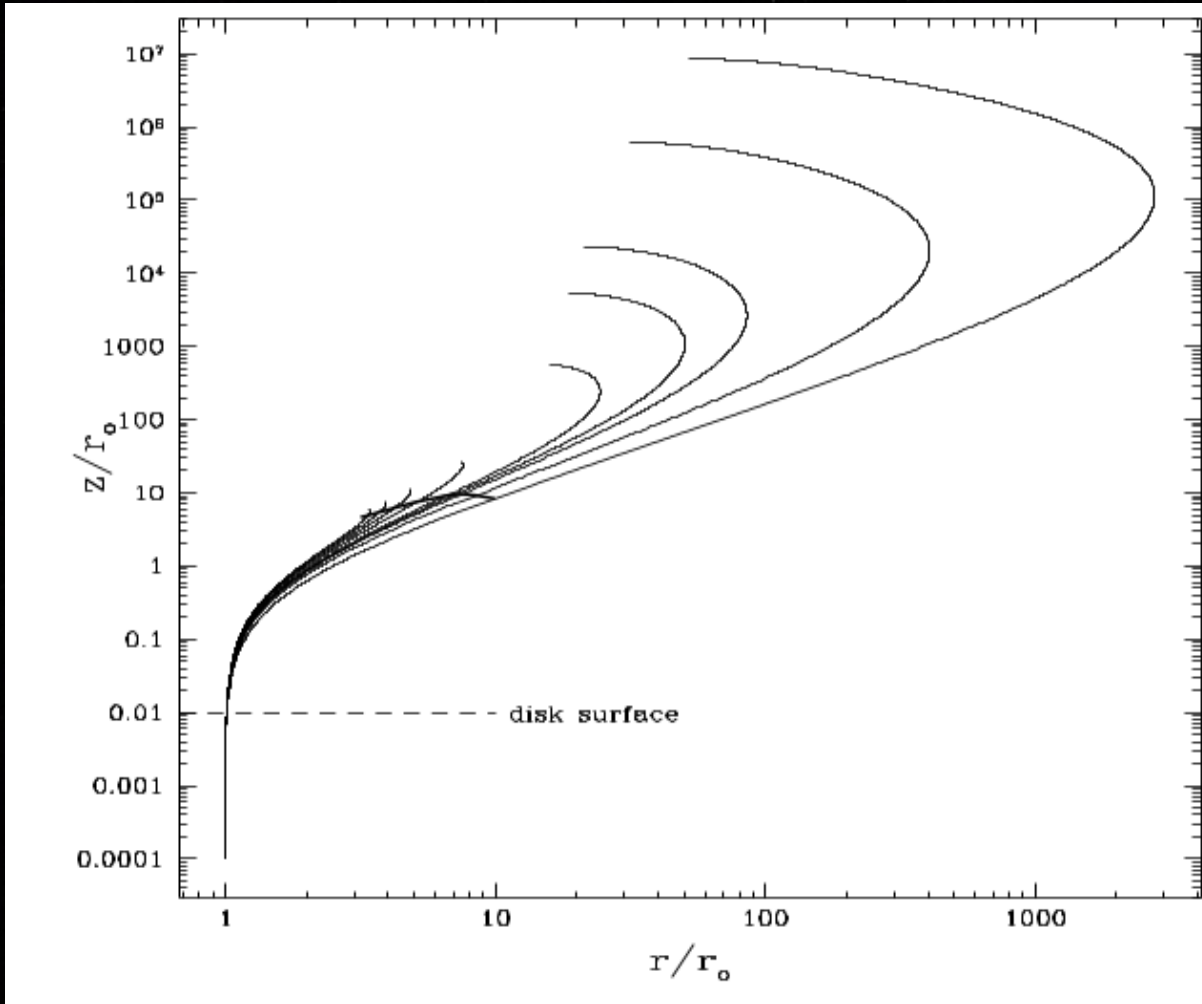
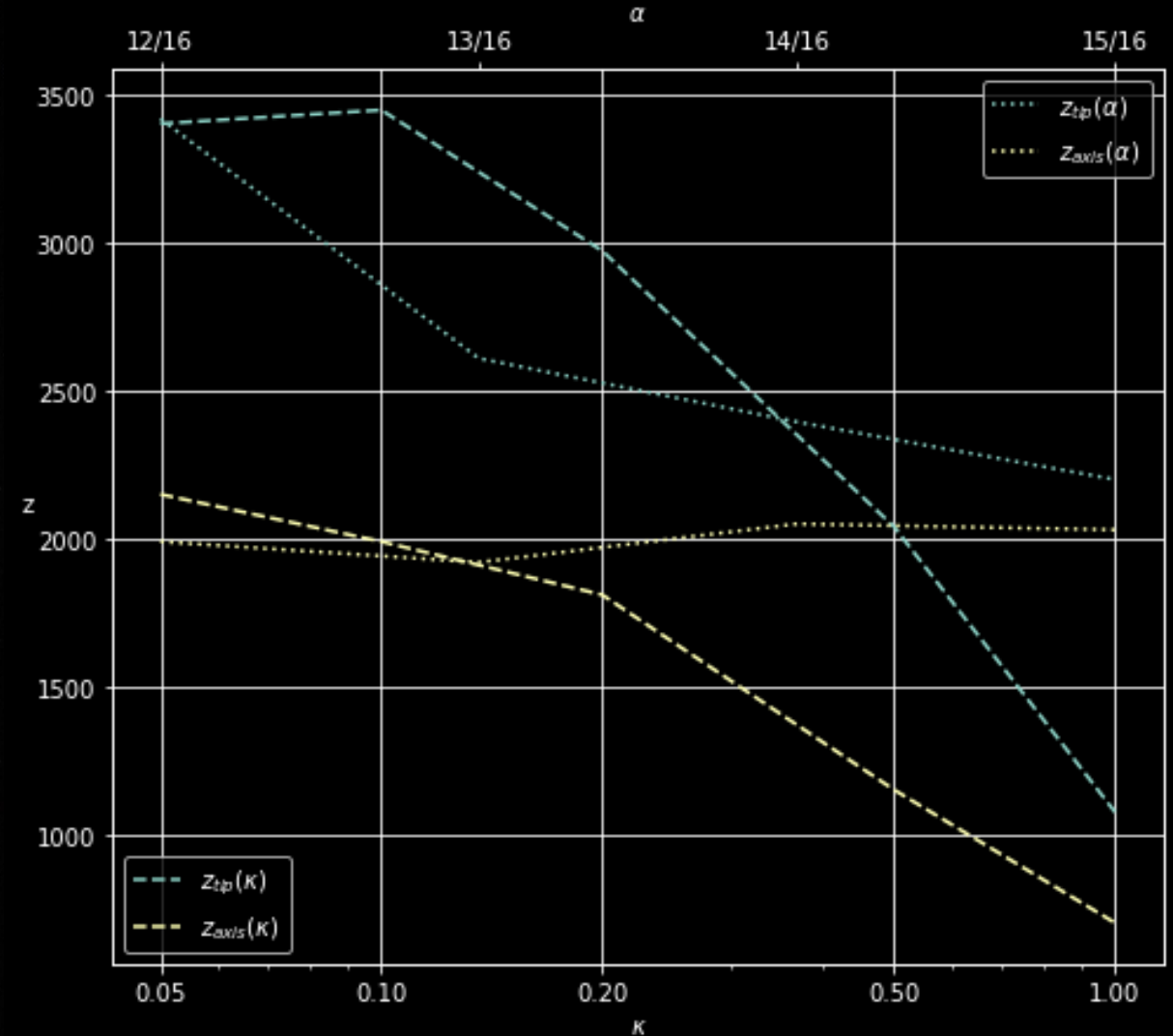
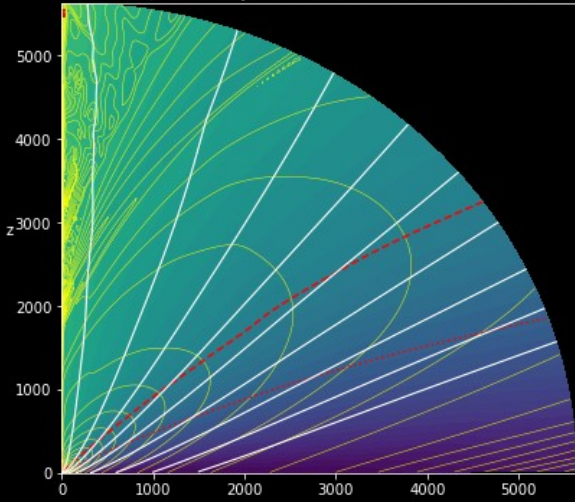


Fig.6 of Ferreira(1997) : magnetic field lines for κ from 0,005 (outermost) à 0,05 (innermost)

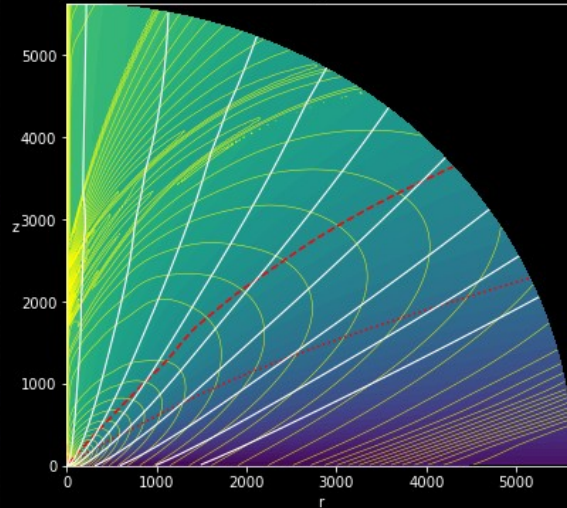


Evolution of the shock altitude Z_{shock} with κ and α

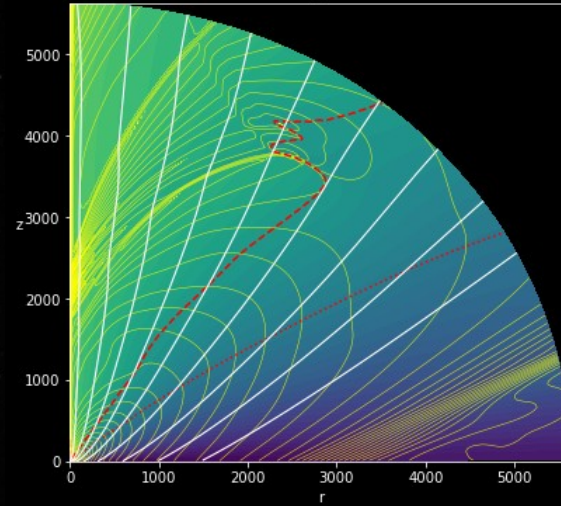
Influence of the magnetic field topology : α



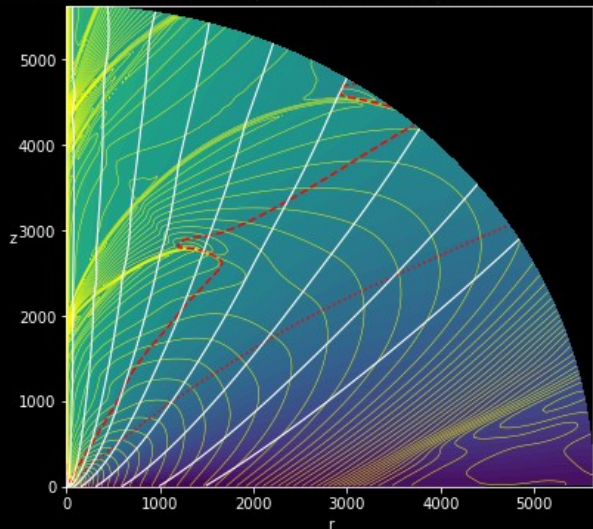
$$\alpha = \frac{10}{16}$$



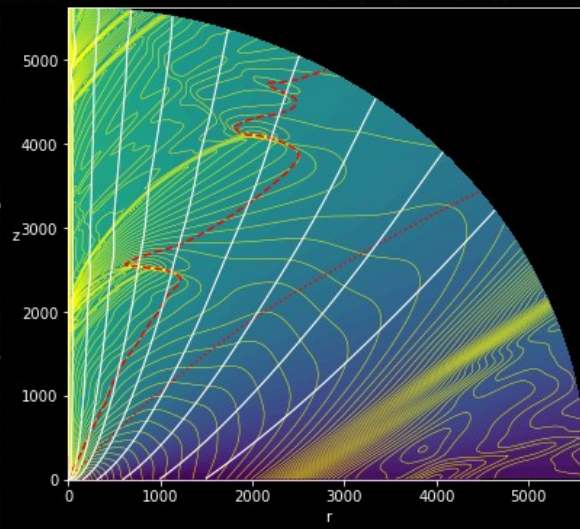
$$\alpha = \frac{11}{16}$$



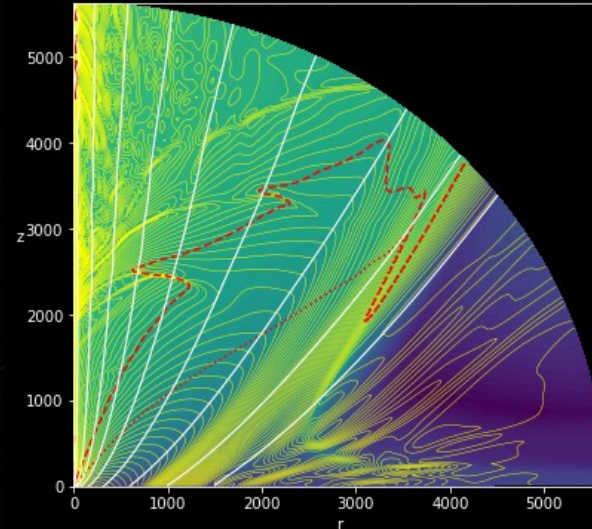
$$\alpha = \frac{12}{16}$$



$$\alpha = \frac{13}{16}$$



$$\alpha = \frac{14}{16}$$



$$\alpha = \frac{15}{16}$$

$$\rho = \rho_d \left(\frac{R}{R_d} \right)^{2\alpha-3}$$

$$P = \rho_d \frac{V_{S_d}^2}{\Gamma} \left(\frac{R}{R_d} \right)^{2\alpha-4}$$

$$B_\theta = -B_{\theta_d} \left(\frac{R}{R_d} \right)^{\alpha-2}$$

$$v_\theta = -v_{\theta_d} \left(\frac{R}{R_d} \right)^{-1/2}$$

$$v_R = v_\theta \frac{B_R}{B_\theta}$$

$$v_\phi = \sqrt{\frac{GM_\star}{R}} + v_\theta \frac{B_\phi}{B_\theta}$$

Influence of the magnetic field topology : α

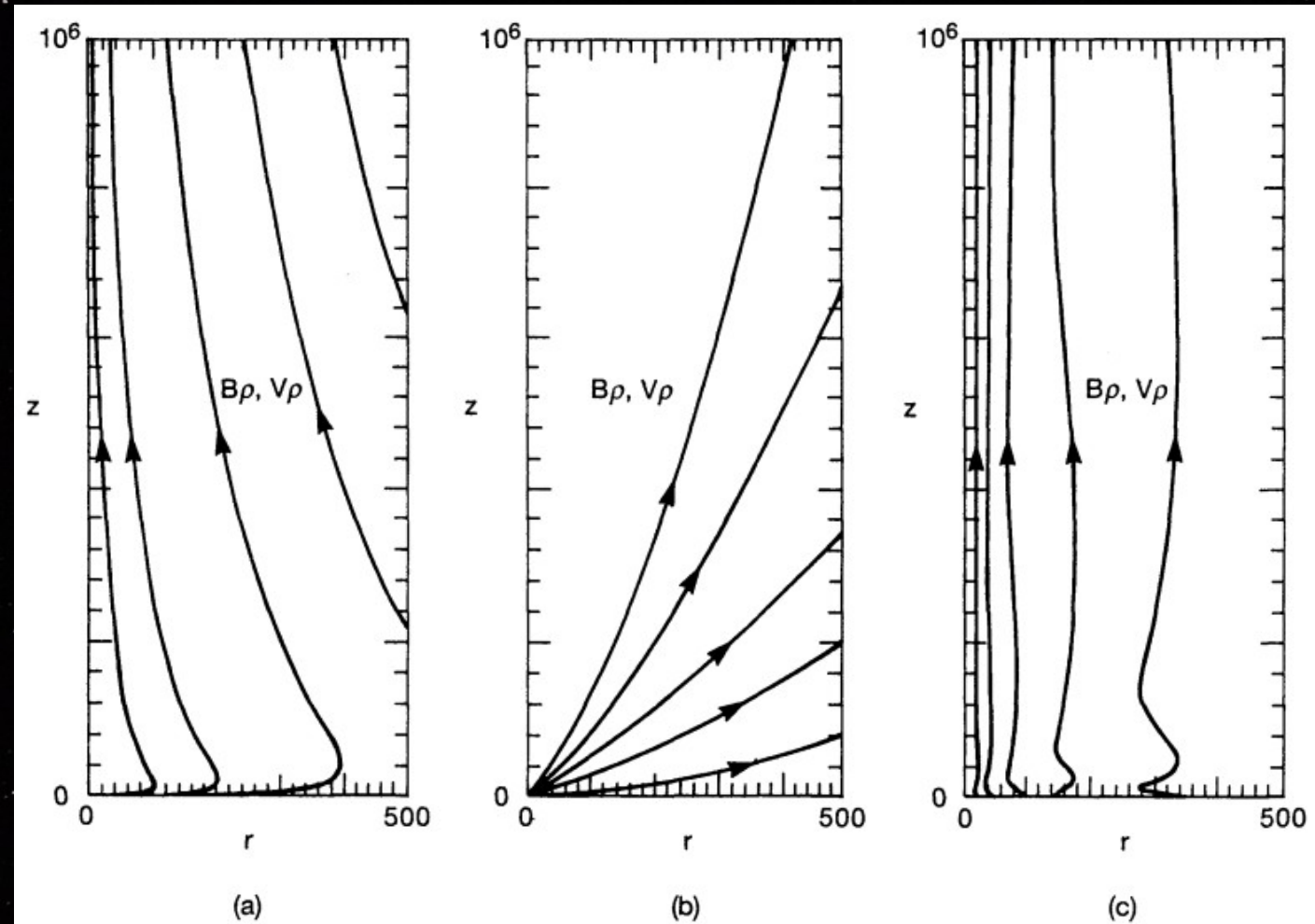
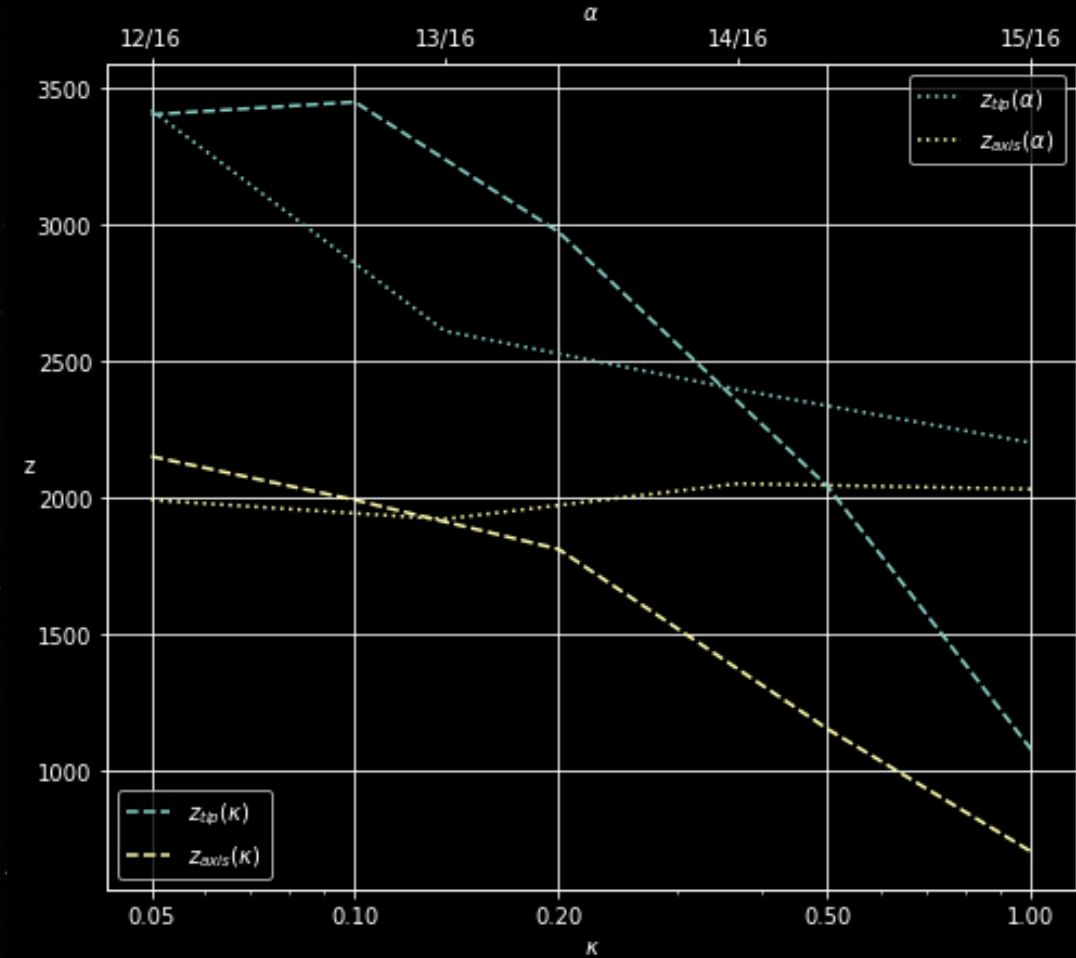
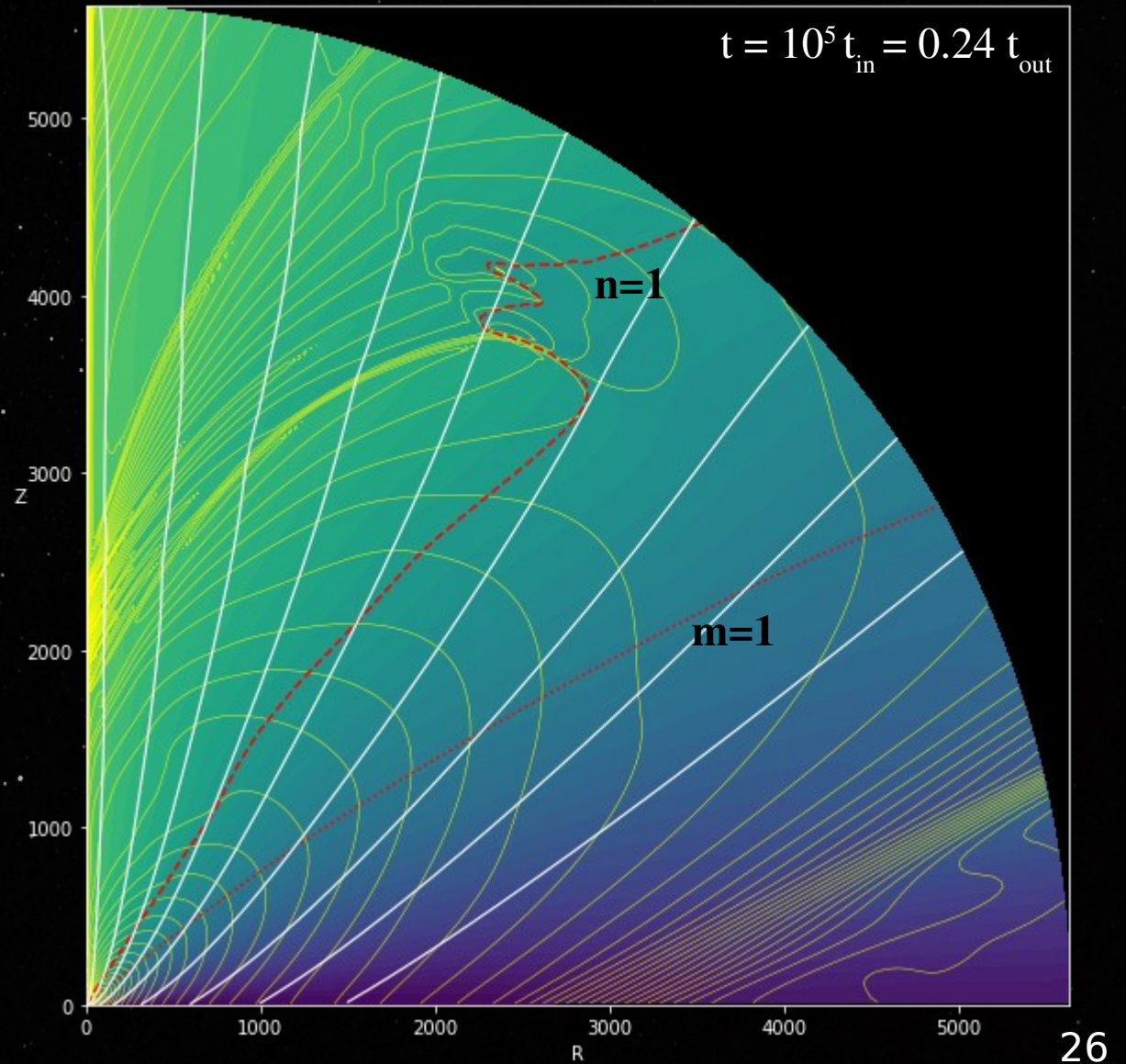
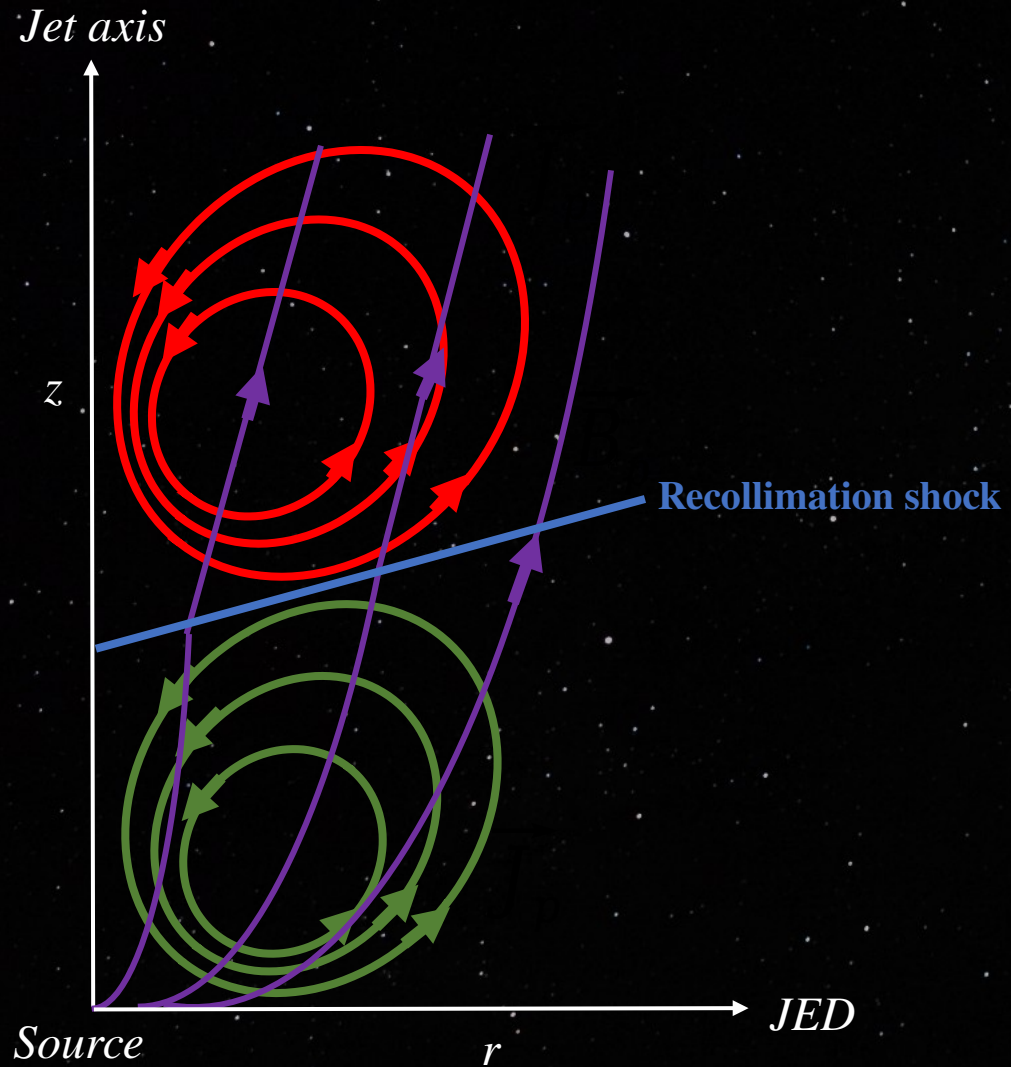


Fig.2 of Contopoulos et Lovelace(1994) : magnetic field lines for $\alpha = 12,8/16$ (a), $\alpha = 14,4/16$ (b) et $\alpha = 16,3/16$ (b)

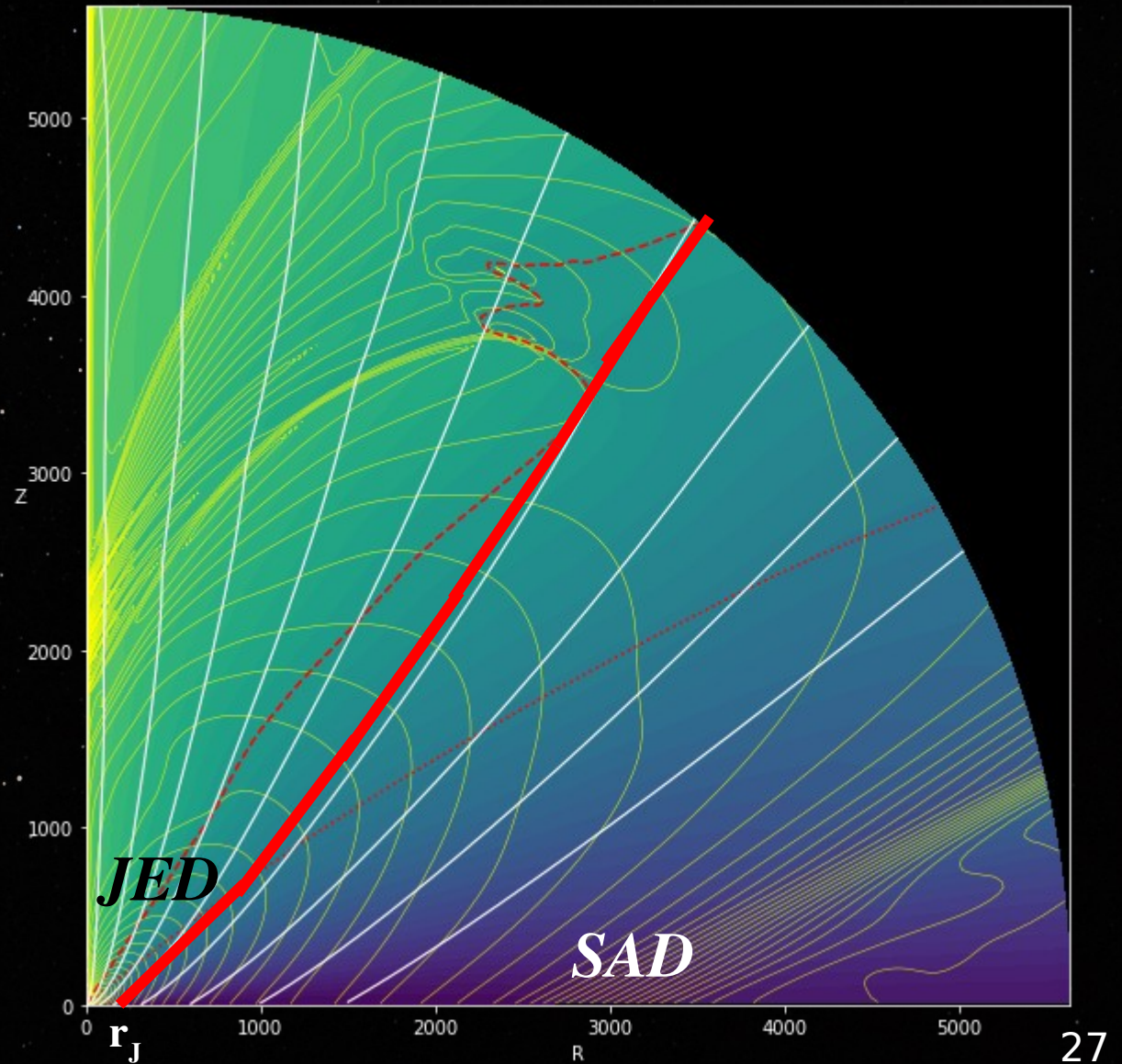


MHD recollimation shocks

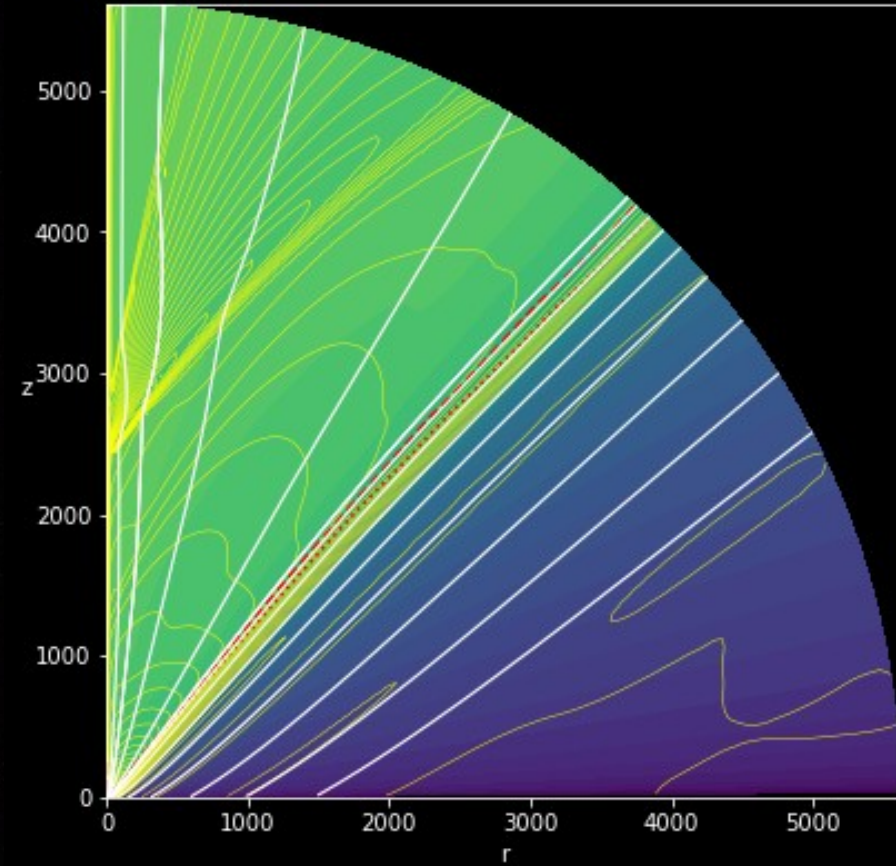


Truncated simulations

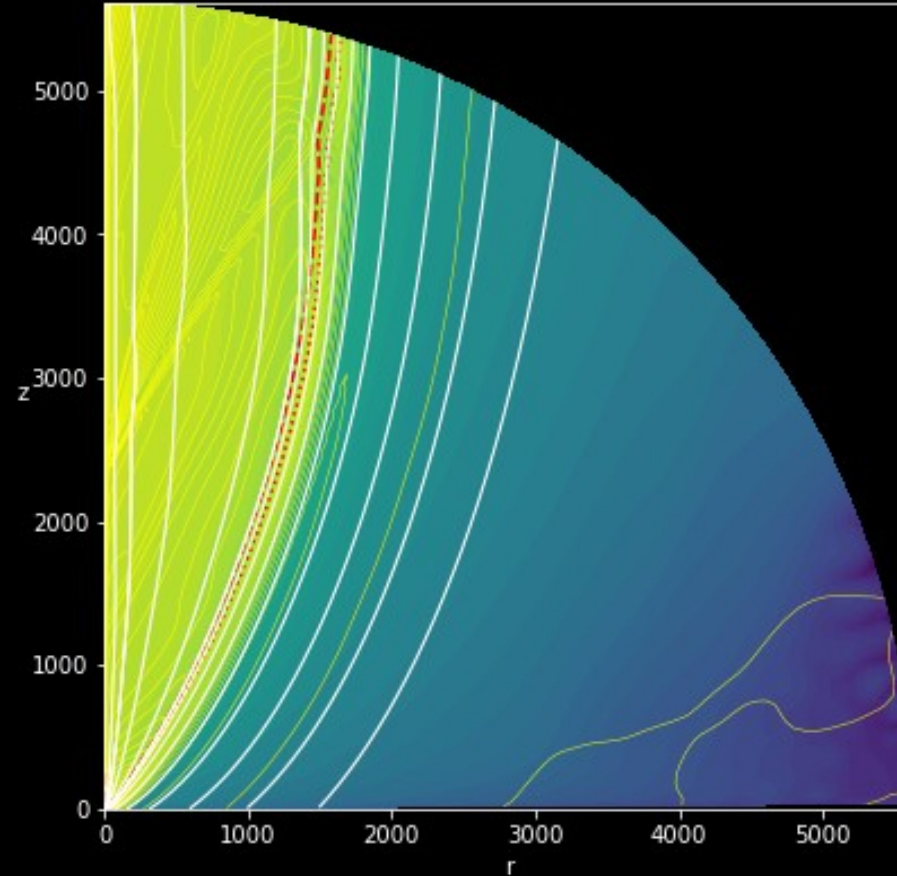
- Self-similar ejection from a finite JED
- No ejection after r_j – Standard Accretion Disk (SAD)
- Observation of shocks ? What characteristics ?



Truncated simulations



No rotation after $R_0=20$



Addition of a constant vertical magnetic field

